

Cosmological Parameter Estimation with Large Scale Structure Observations

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in collaboration with

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arXiv:1307.1459

arXiv:1308.6186

Heidelberg

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Center for Astroparticle Physics
GENEVA



UNIVERSITÉ
DE GENÈVE

Outline

Introduction

Galaxy number counts

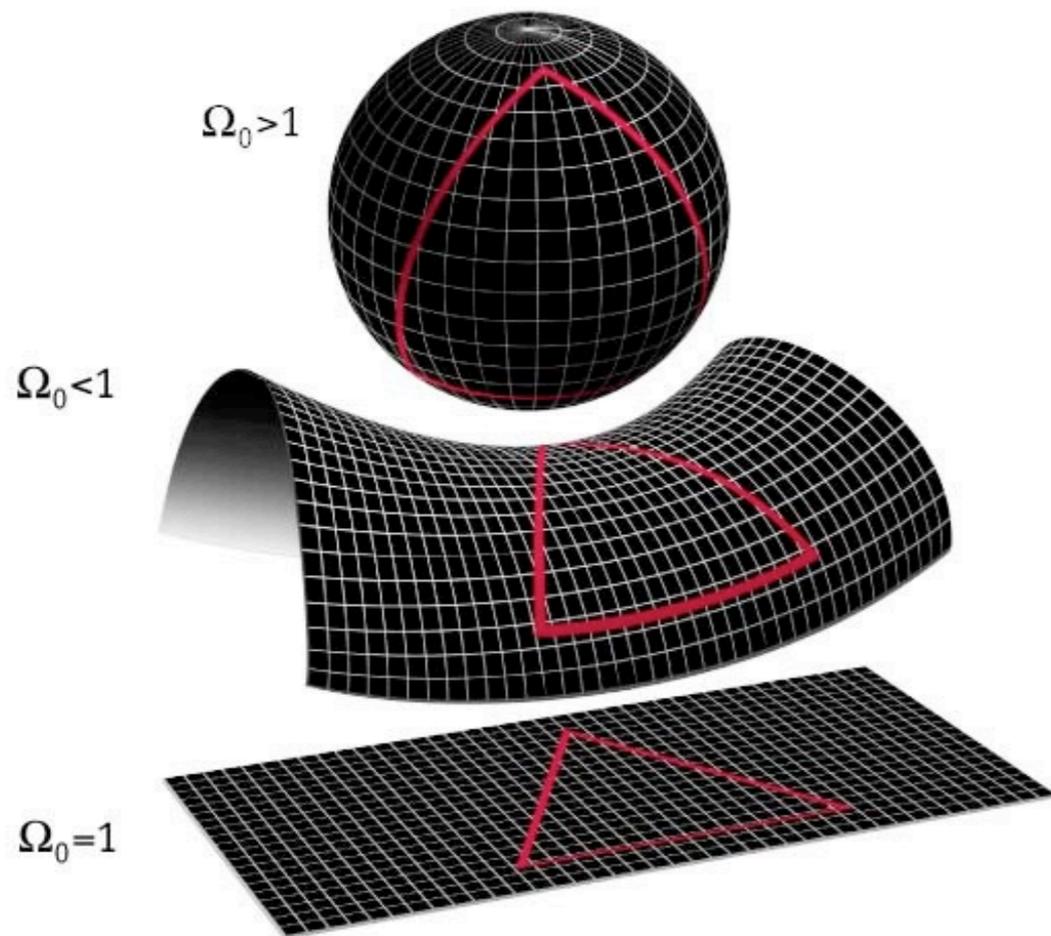
2D redshift dependent power spectra

Fisher matrix analysis

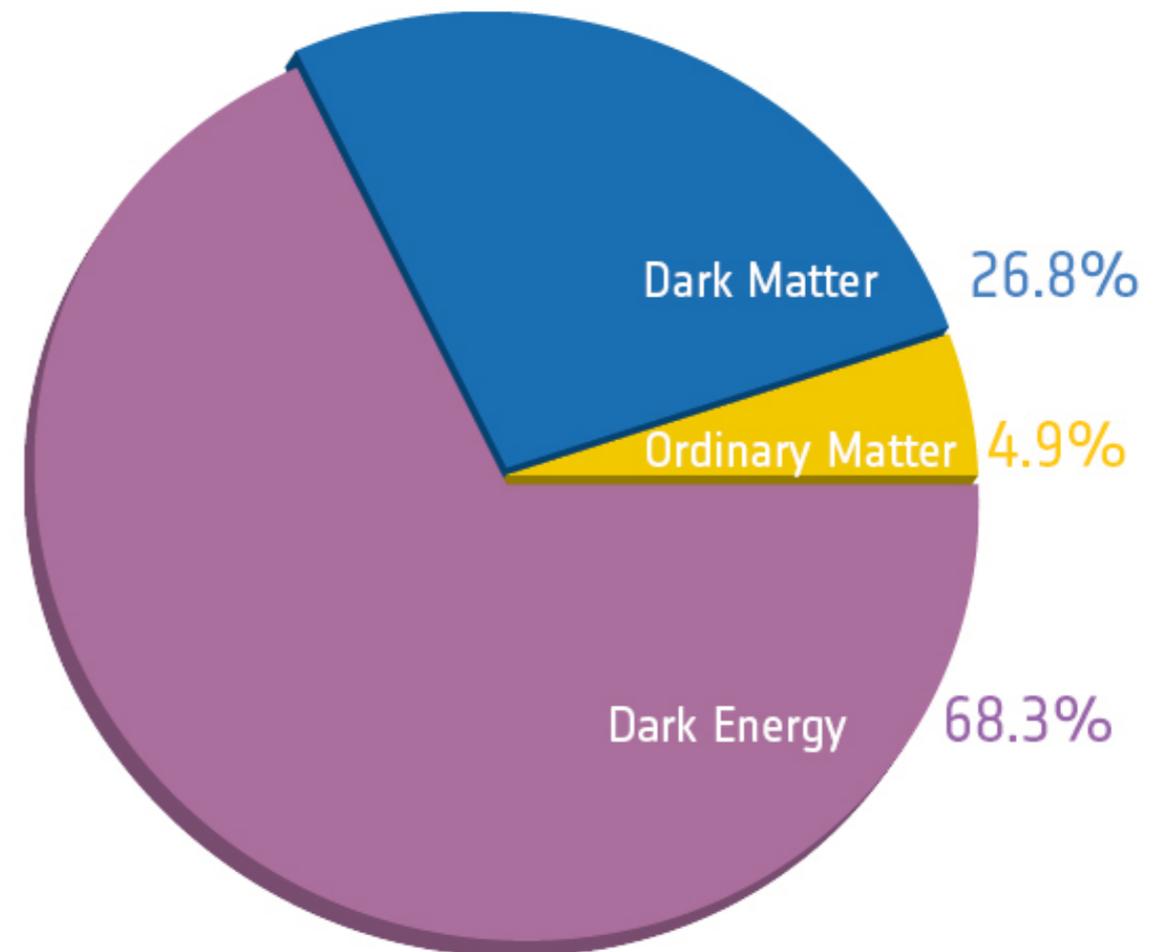
Lensing Potential

CLASSgal

Standard Cosmological Model

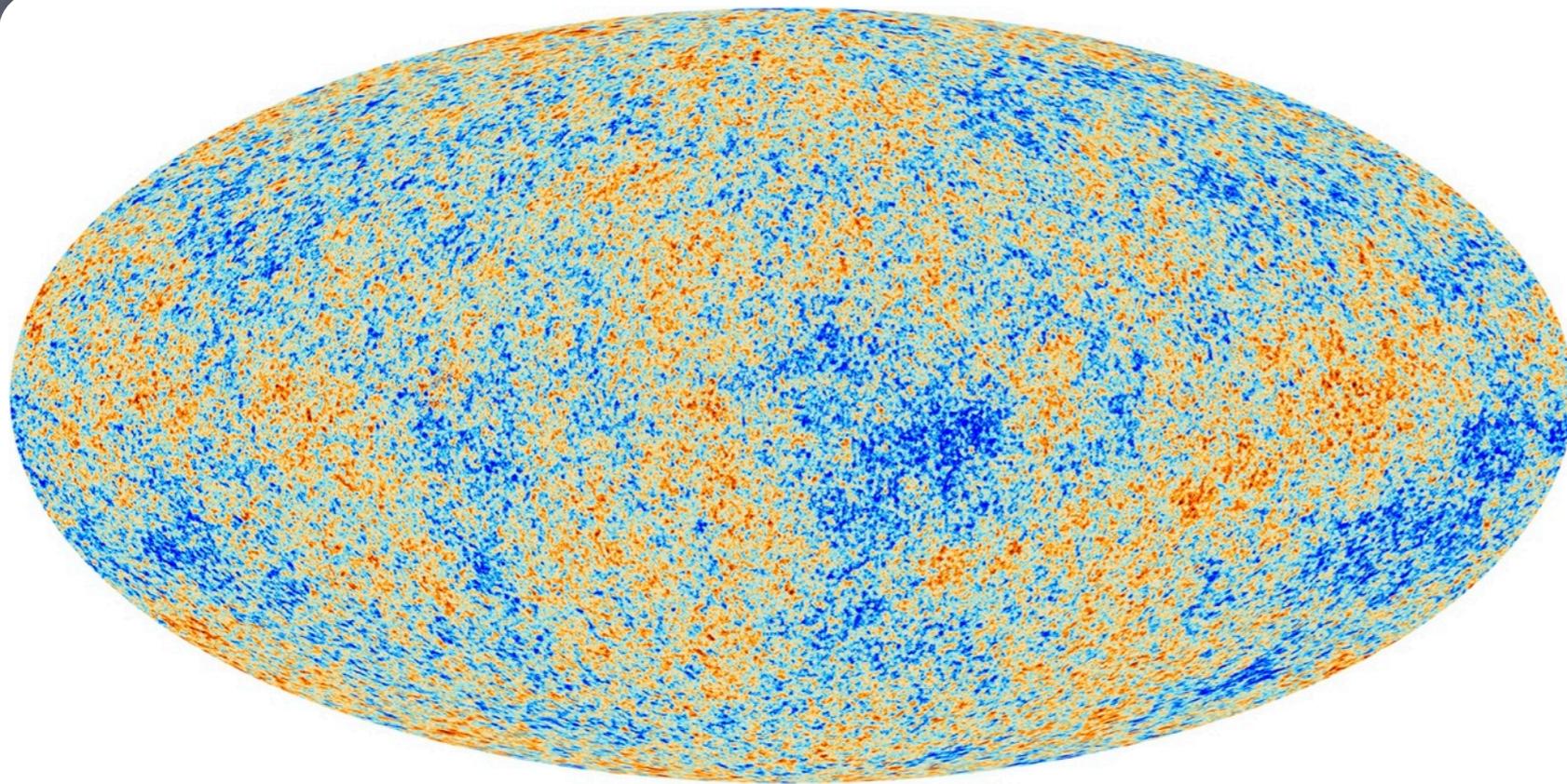


MAP990006

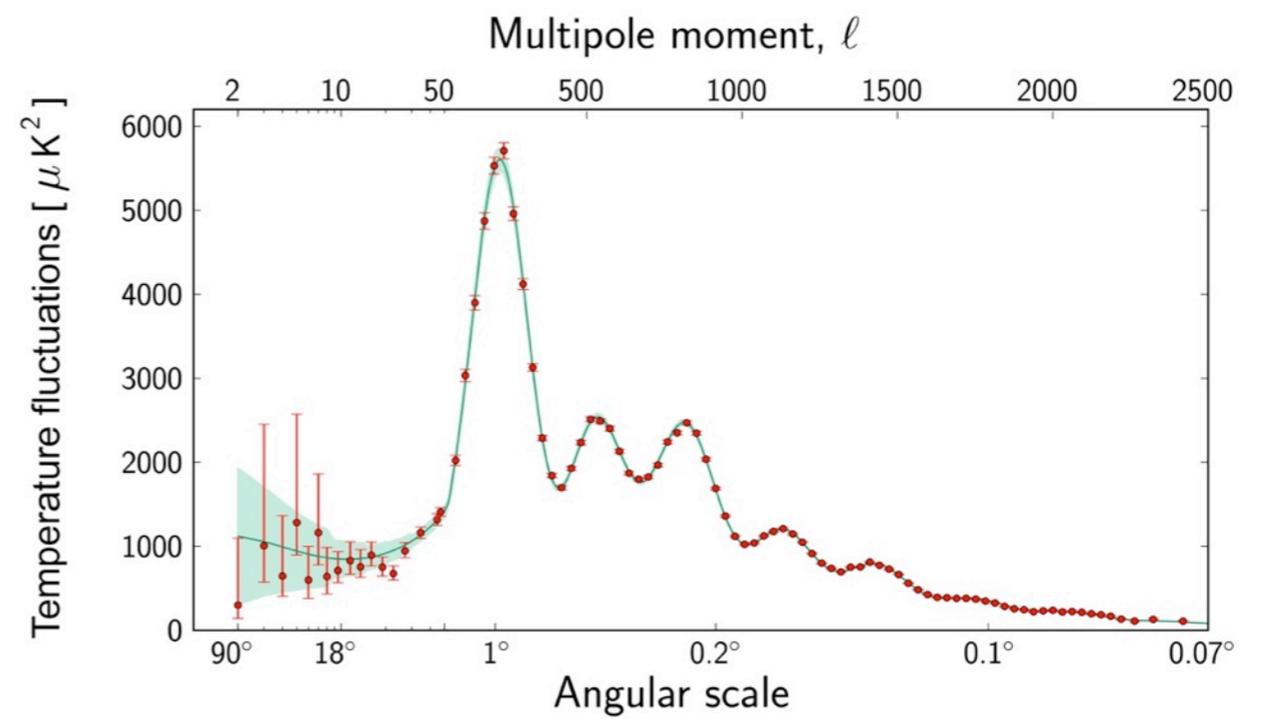


$$G_{\mu\nu} = \kappa T_{\mu\nu}$$

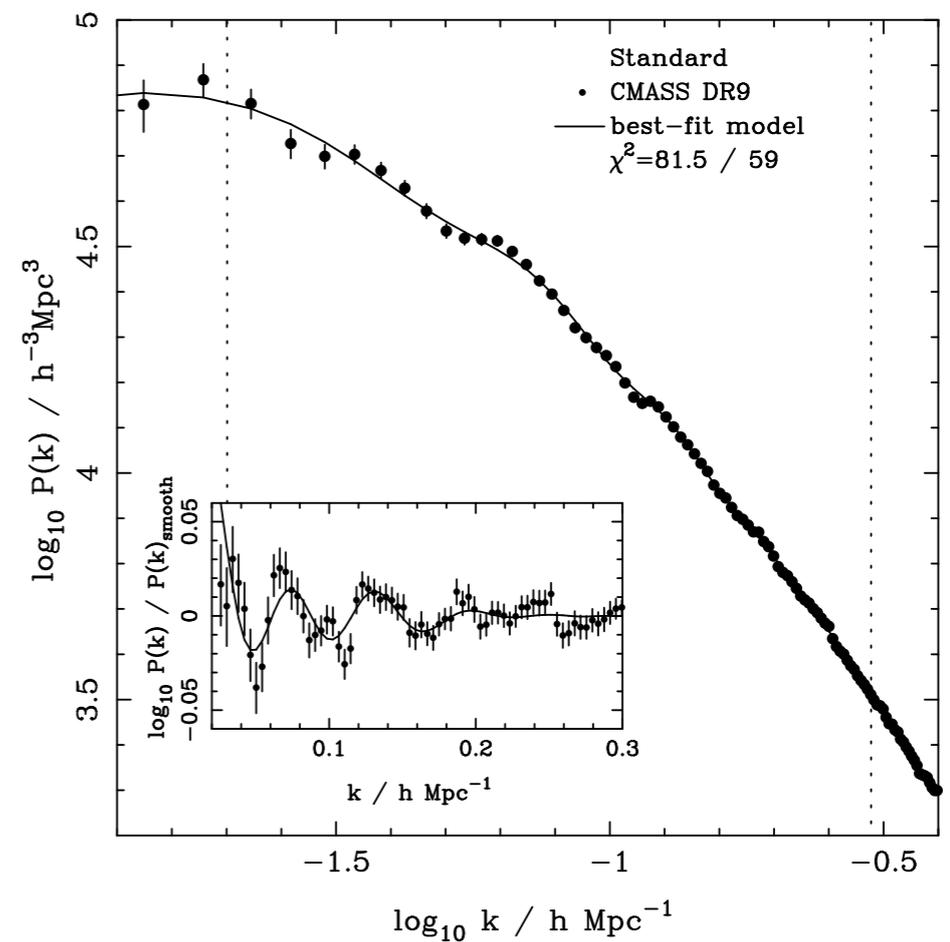
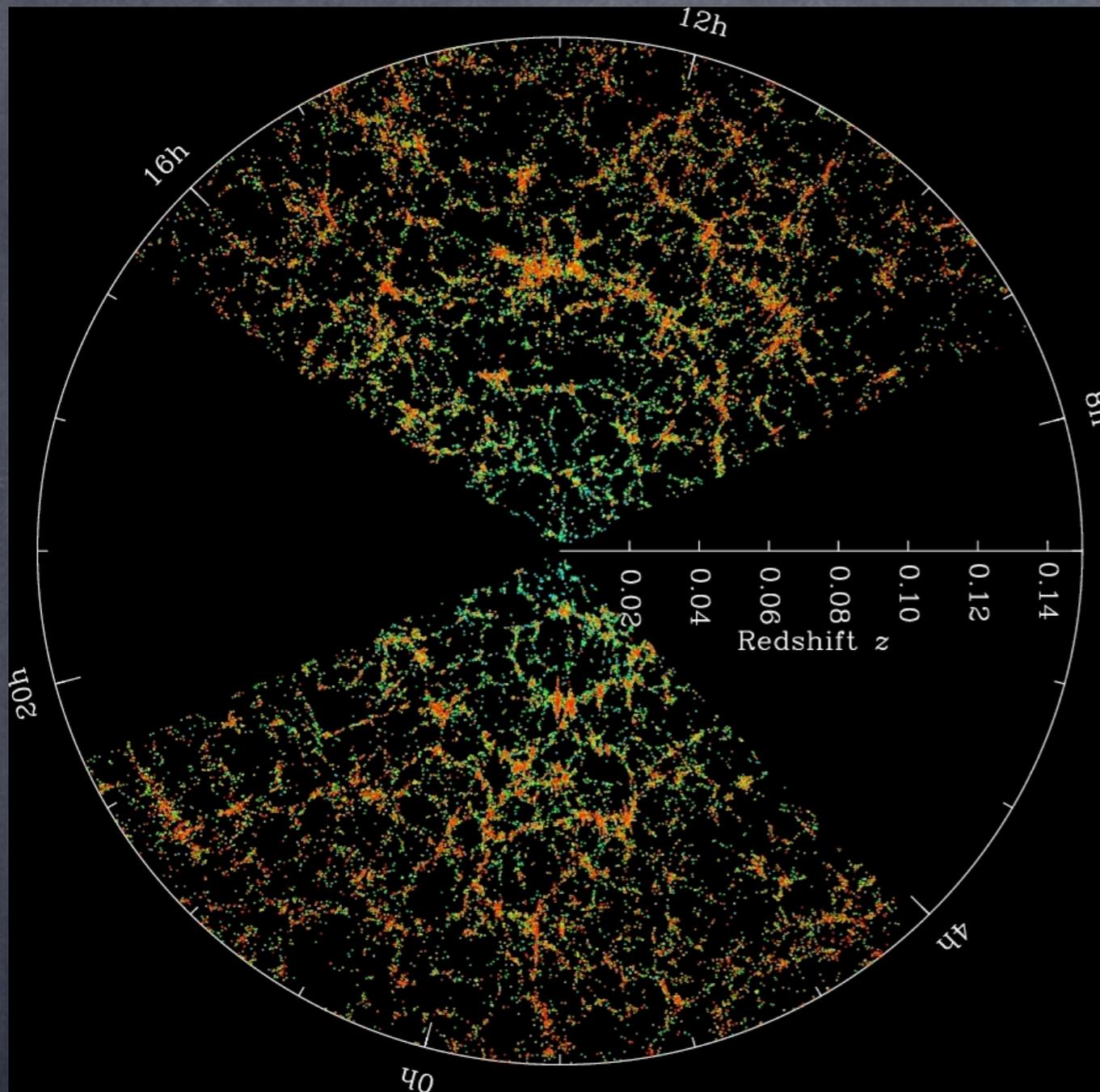
CMB



Planck Collaboration

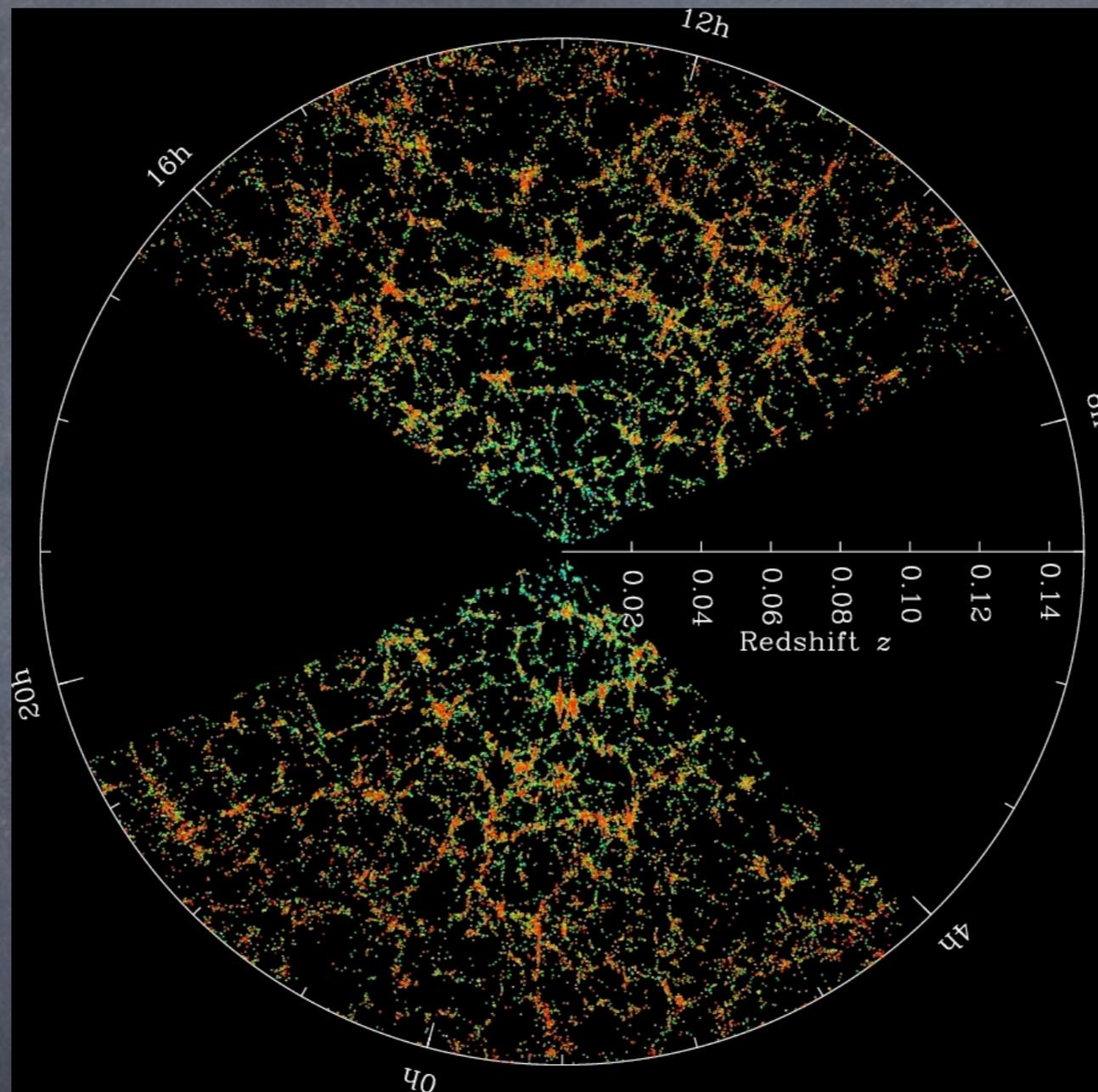


Large Scale Structures



Anderson et al '12 [arXiv:1203.6594]

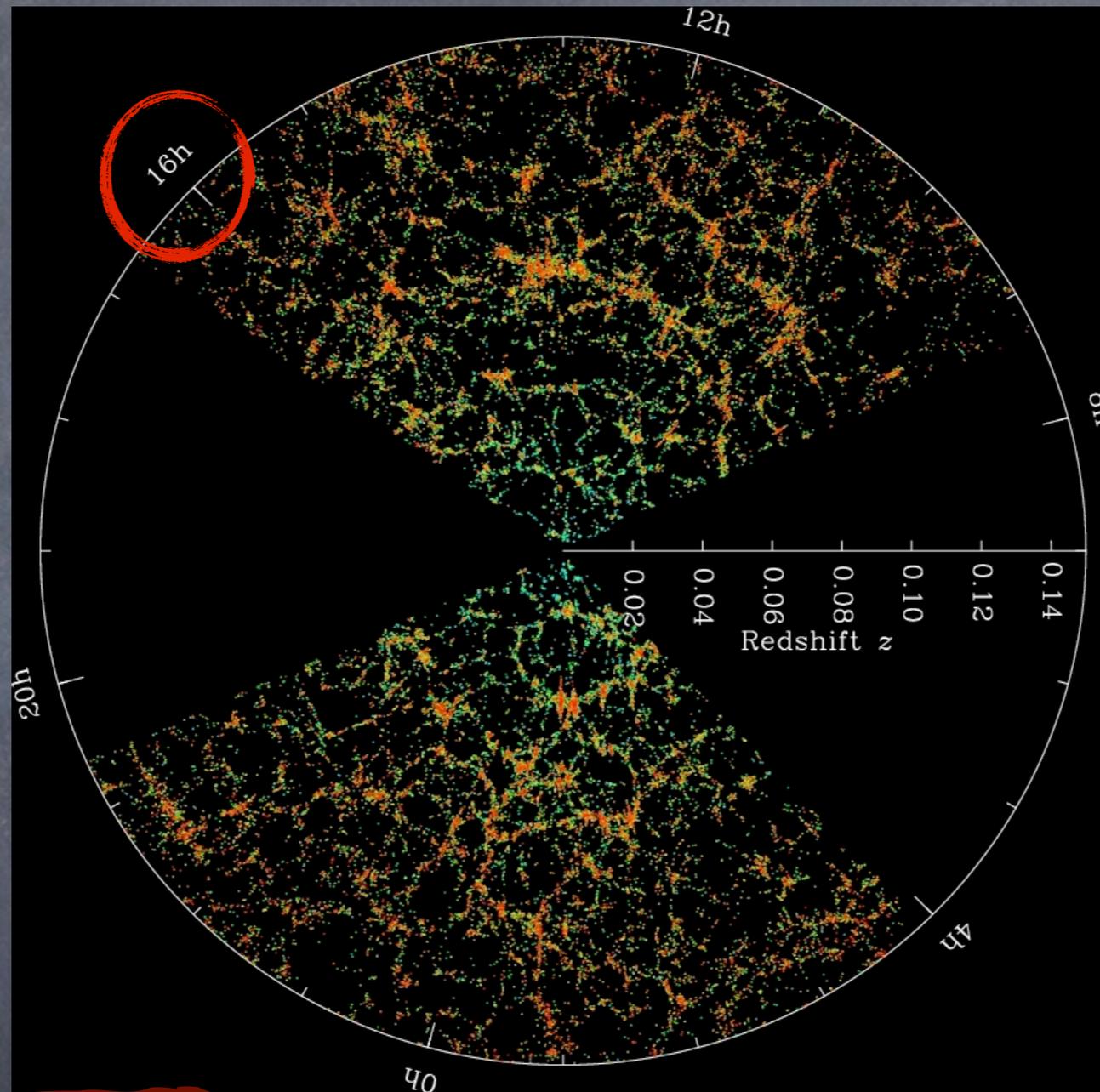
What do we really observe?



Angular position \mathbf{n} Redshift z

$$N(\mathbf{n}, z) d\Omega_{\mathbf{n}} dz$$

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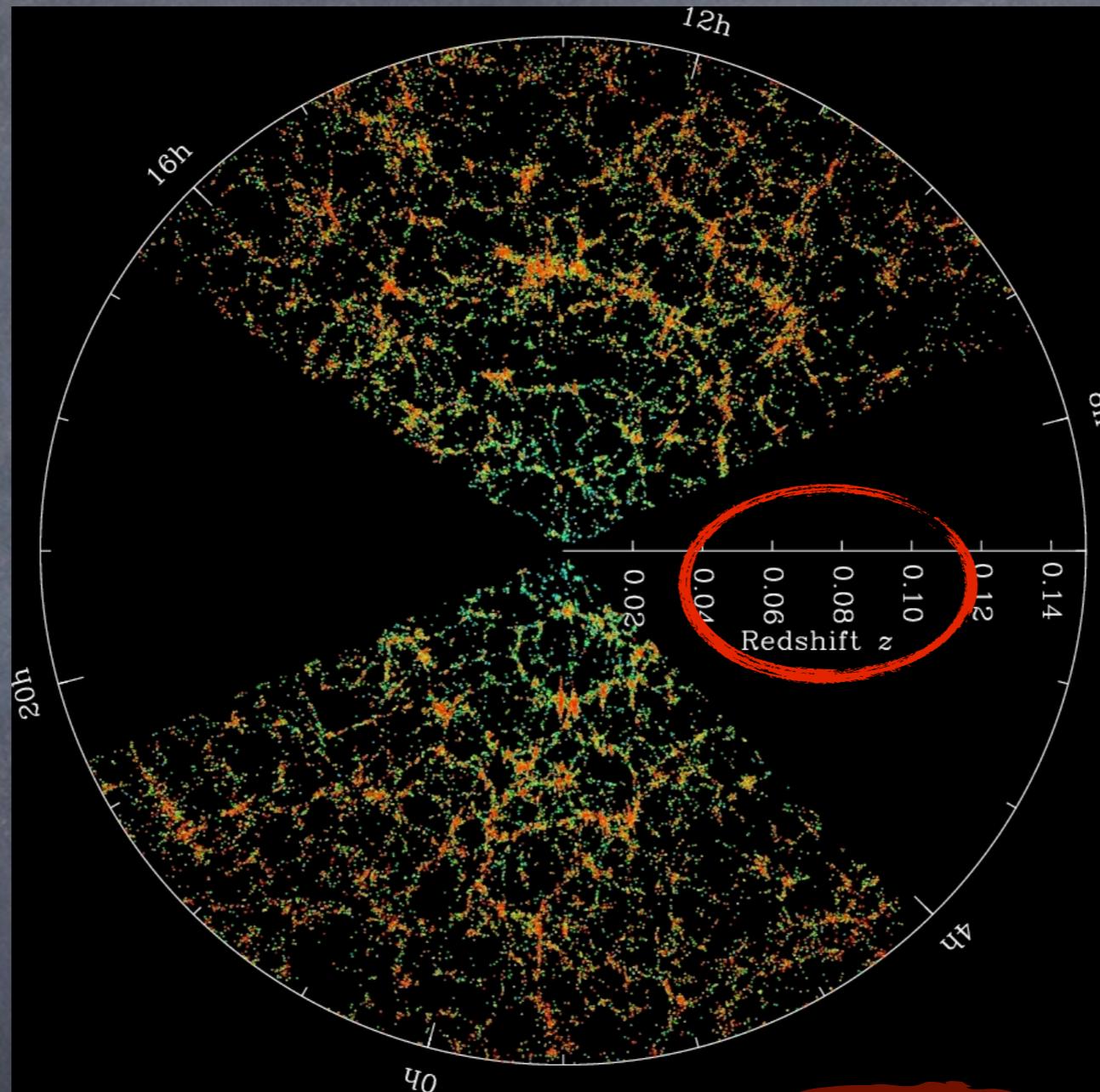


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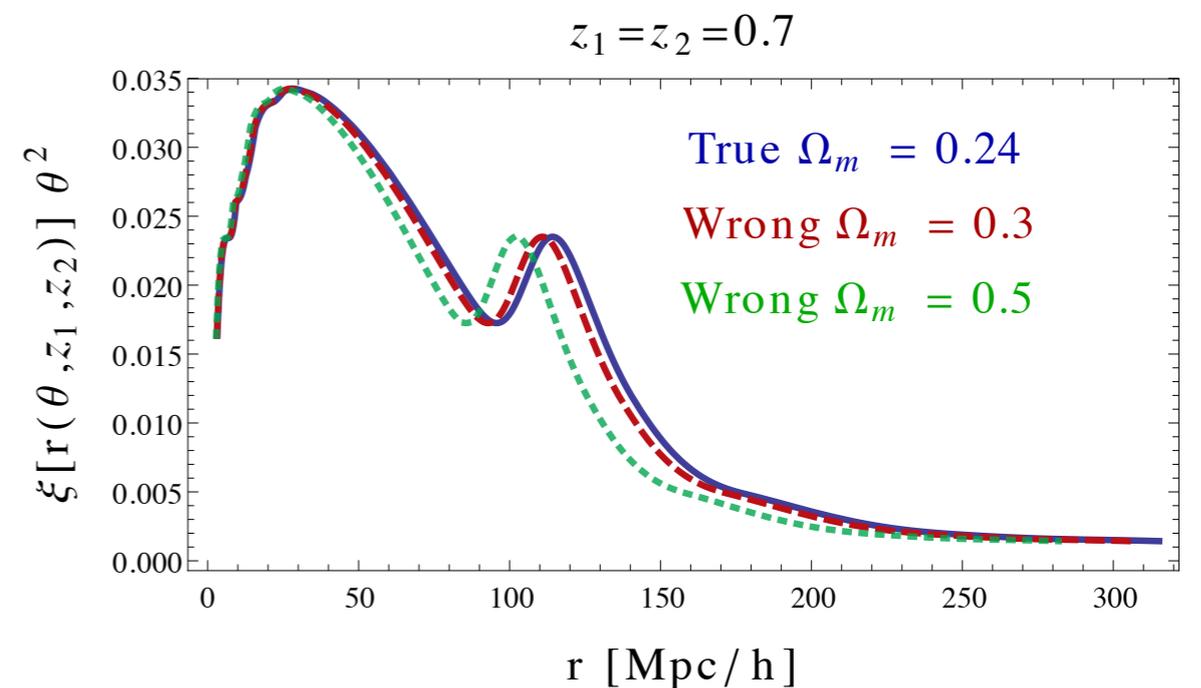
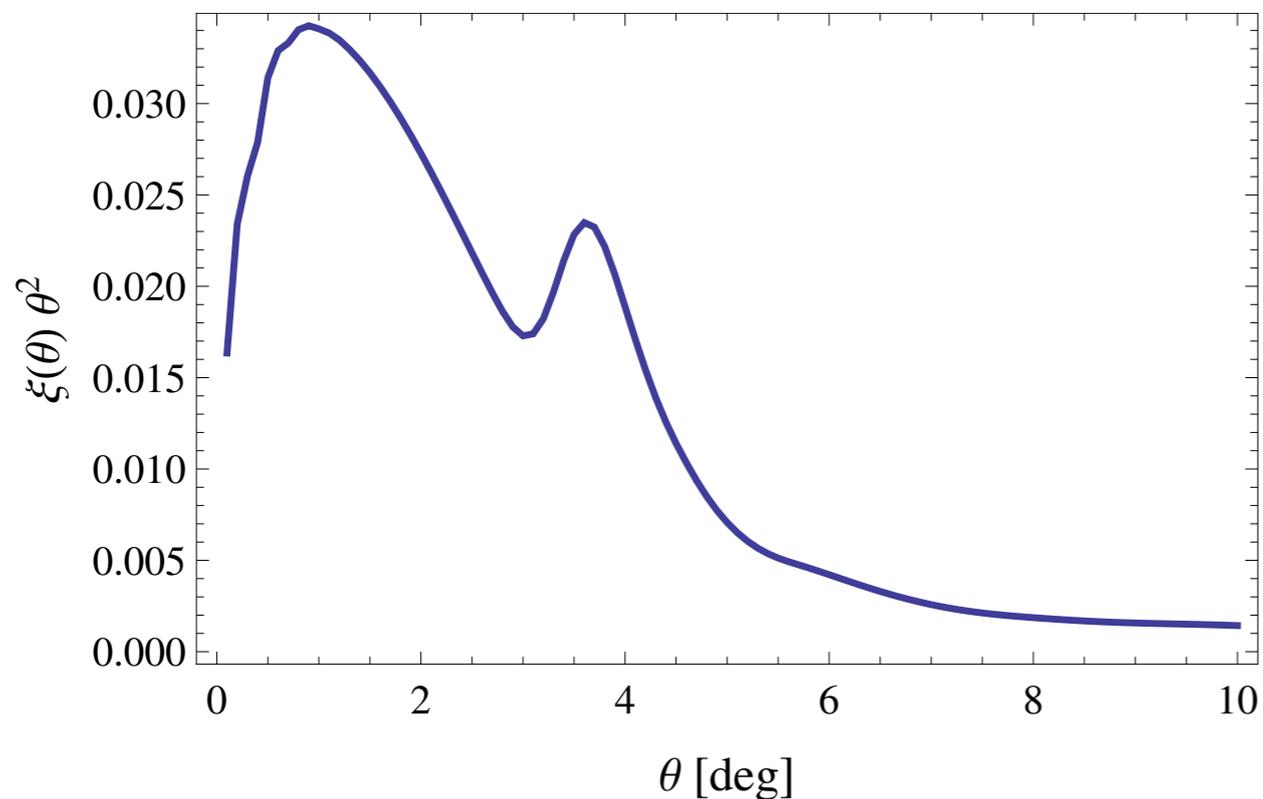
$$N(\mathbf{n}, z) d\Omega_{\mathbf{n}} dz$$

What do we really observe?

In 3D analysis we need to convert a redshift and angles into length scales

$$r(z_1, z_2, \theta) = \sqrt{\chi^2(z_1) + \chi^2(z_2) - 2\chi(z_1)\chi(z_2) \cos \theta}$$

We need to assume a cosmology!

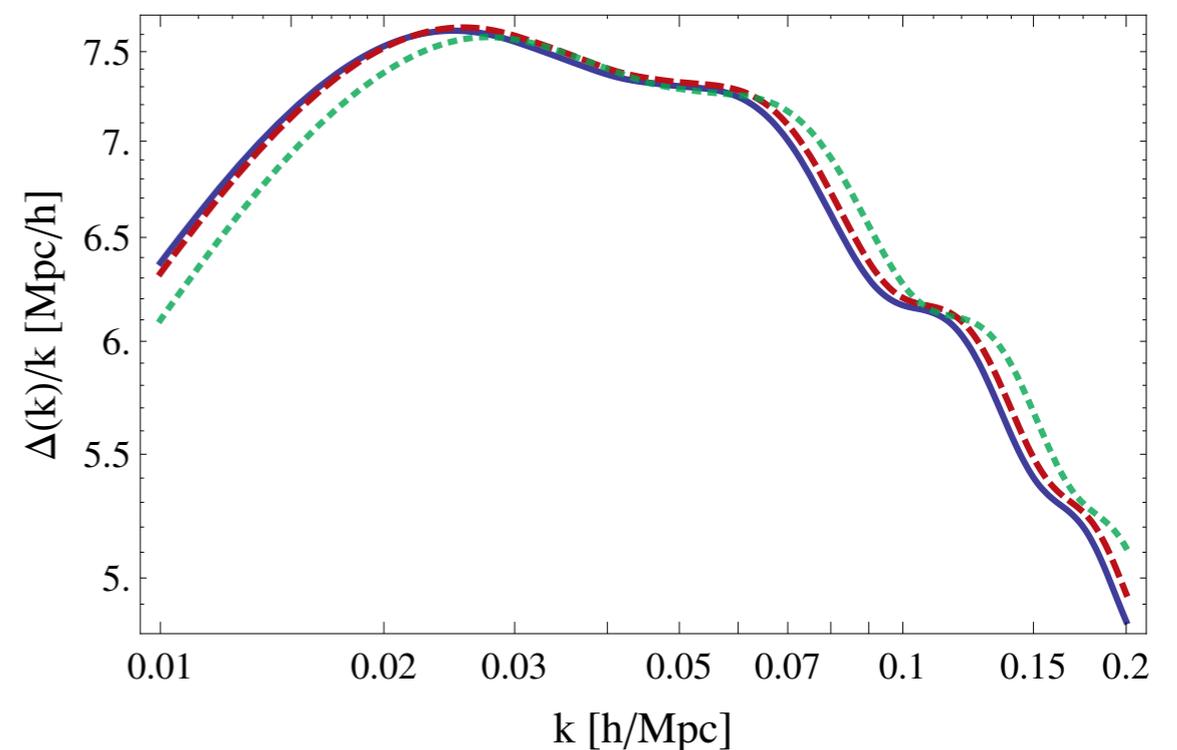
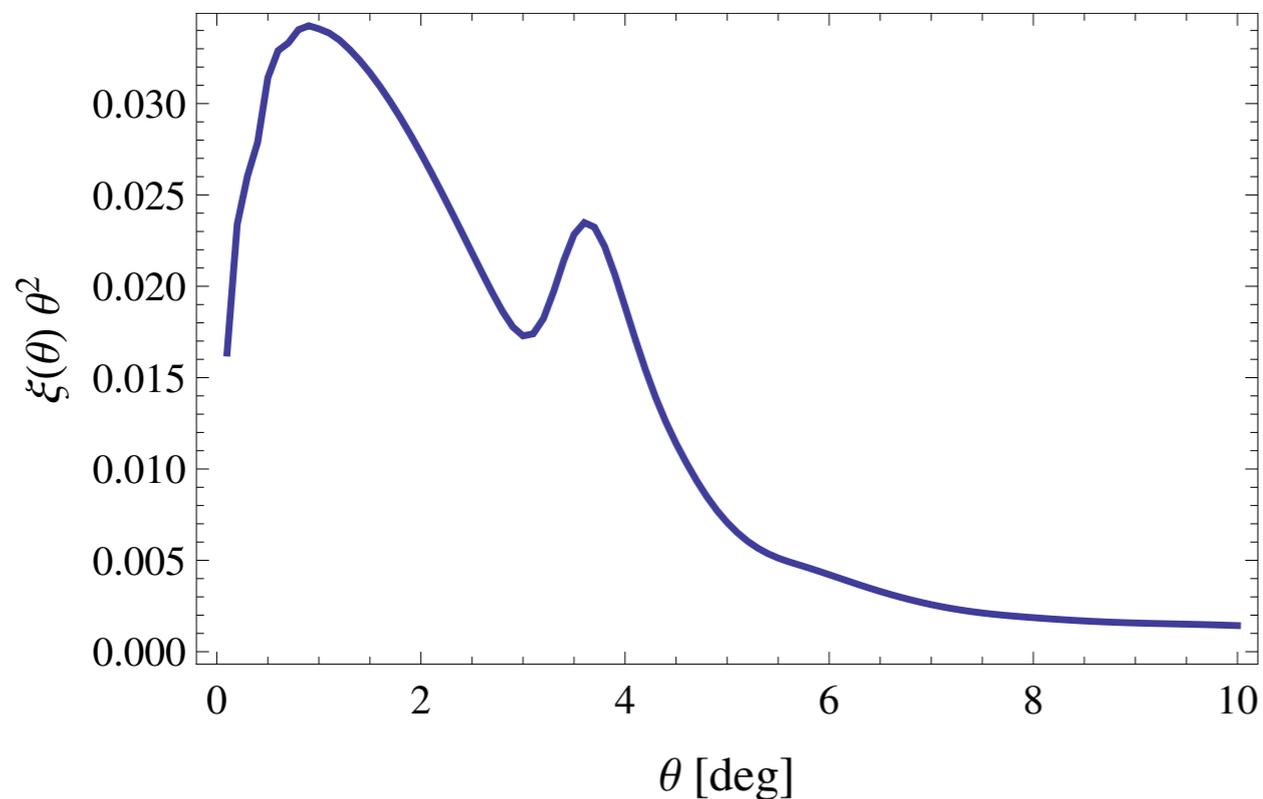


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What do we really observe?

To compute $\Delta(\mathbf{n}, z) \equiv \frac{N(\mathbf{n}, z) - \langle N \rangle(z)}{\langle N \rangle(z)}$
we have to consider:

- observation on the past lightcone
- redshift perturbed by peculiar velocity
- light deflection
- volume distortion

$$\begin{aligned}\Delta(\mathbf{n}, z) = & D_g + \Phi + \Psi + \frac{1}{\mathcal{H}} [\Phi' + \partial_r(\mathbf{V} \cdot \mathbf{n})] \\ & + \left(\frac{\mathcal{H}'}{\mathcal{H}^2} + \frac{2}{r_S \mathcal{H}} \right) \left(\Psi + \mathbf{V} \cdot \mathbf{n} + \int_0^{r_S} dr (\Phi' + \Psi') \right) \\ & + \frac{1}{r_S} \int_0^{r_S} dr \left[2 - \frac{r_S - r}{r} \Delta_\Omega \right] (\Phi + \Psi)\end{aligned}$$

Bonvin & Durrer [arXiv:1105.5280],
Challinor & Lewis [arXiv:1105.5292],
Yoo [arXiv:1009.3021]

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What

To compare
we have

- observed
- redshift
- light
- volume

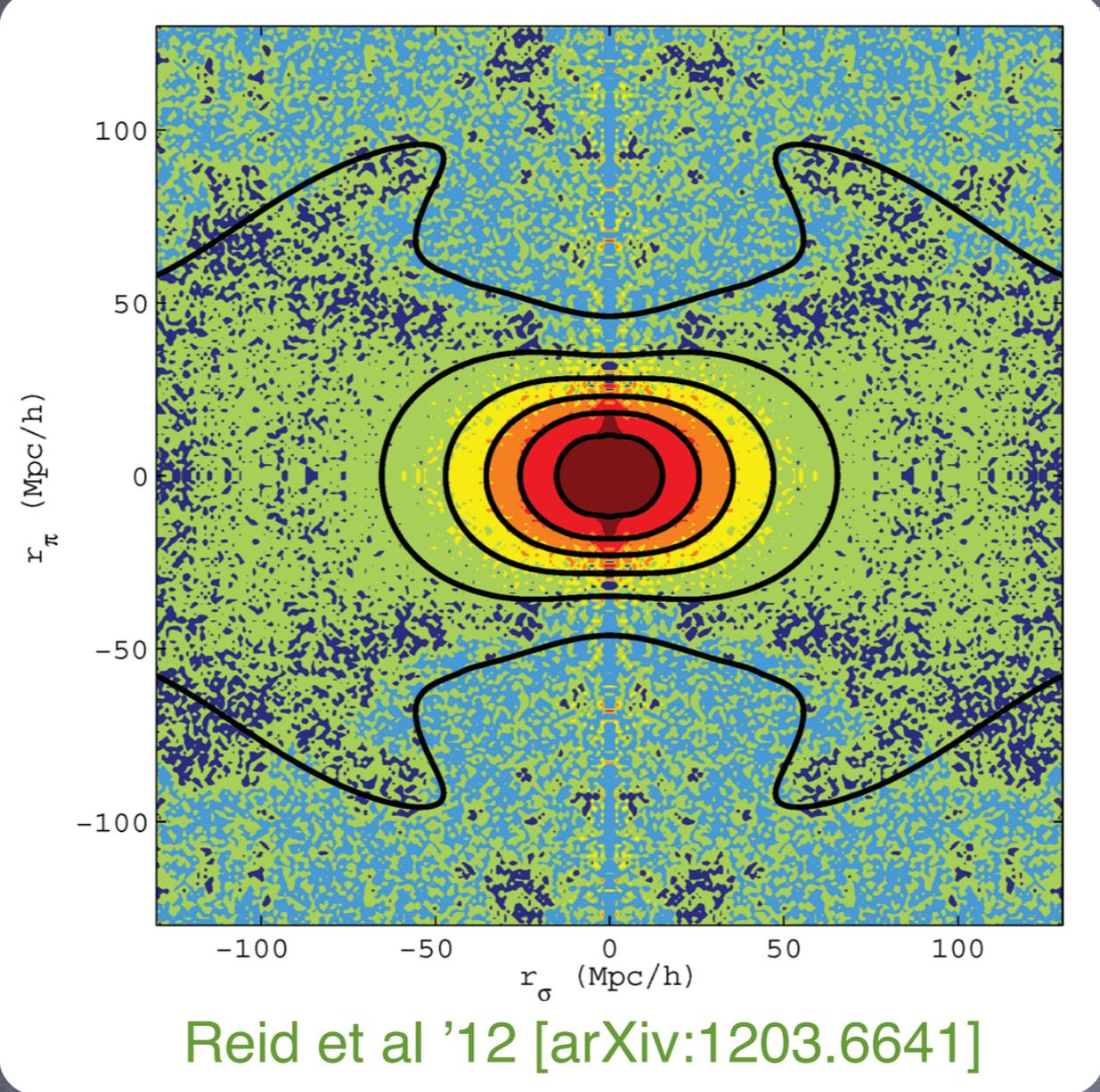
curve?

$\gamma > (z)$

(z)

e

velocity



$$\Delta(\mathbf{n}, z) = D_g + \Phi + \Psi + \frac{1}{\mathcal{H}} [\Phi' + \partial_r(\mathbf{V} \cdot \mathbf{n})] + \left(\frac{\mathcal{H}'}{\mathcal{H}^2} + \frac{2}{r_S \mathcal{H}} \right) \left(\Psi + \mathbf{V} \cdot \mathbf{n} + \int_0^{r_S} dr (\Phi' + \Psi') \right) + \frac{1}{r_S} \int_0^{r_S} dr \left[2 - \frac{r_S - r}{r} \Delta_\Omega \right] (\Phi + \Psi)$$

z-distortion

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Bonvin & Durrer [arXiv:1105.5280],
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What do we really observe?

Additional effects

galaxy bias

magnification bias

galaxy evolution

$$\begin{aligned}\Delta^{(N)}(\mathbf{n}, z, m_*) &= bD_g(L > \bar{L}_*) + (1 + 5s)\Phi + \Psi + \frac{1}{\mathcal{H}} [\Phi' + \partial_r(\mathbf{V} \cdot \mathbf{n})] \\ &+ \left(\frac{\mathcal{H}'}{\mathcal{H}^2} + \frac{2 - 5s}{r_S \mathcal{H}} + 5s - f_{\text{evo}}^N \right) \left(\Psi + \mathbf{V} \cdot \mathbf{n} + \int_0^{r_S} dr (\Phi' + \Psi') \right) \\ &+ \frac{2 - 5s}{2r_S} \int_0^{r_S} dr \left[2 - \frac{r_S - r}{r} \Delta_\Omega \right] (\Phi + \Psi)\end{aligned}$$

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$$s(z, m_*) \equiv \frac{\partial \log_{10} \bar{N}(z, m < m_*)}{\partial m_*},$$

$$\bar{N}(z, L_S < \bar{L}_*) = \int_{F_*}^{\infty} \bar{n}_S(z, \ln F) d \ln F$$

$$\begin{aligned} \Delta^{(N)}(\mathbf{n}, z, m_*) &= bD_g(L > \bar{L}_*) + (1 + 5s)\Phi + \Psi + \frac{1}{\mathcal{H}} [\Phi' + \partial_r(\mathbf{V} \cdot \mathbf{n})] \\ &+ \left(\frac{\mathcal{H}'}{\mathcal{H}^2} + \frac{2 - 5s}{r_S \mathcal{H}} - 5s - f_{\text{evo}}^N \right) \left(\Psi + \mathbf{V} \cdot \mathbf{n} + \int_0^{r_S} dr (\Phi' + \Psi') \right) \\ &+ \frac{2 - 5s}{2r_S} \int_0^{r_S} dr \left[2 - \frac{r_S - r}{r} \Delta_\Omega \right] (\Phi + \Psi) \end{aligned}$$

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galaxy evolution

$$f_{\text{evo}}^N = \frac{\partial \ln(a^3 \bar{N}(z, L > \bar{L}_*))}{\mathcal{H} \partial \tau_S}$$

$$\begin{aligned} \Delta^{(N)}(\mathbf{n}, z, m_*) &= bD_g(L > \bar{L}_*) + (1 + 5s)\Phi + \Psi + \frac{1}{\mathcal{H}} [\Phi' + \partial_r(\mathbf{V} \cdot \mathbf{n})] \\ &+ \left(\frac{\mathcal{H}'}{\mathcal{H}^2} + \frac{2 - 5s}{r_S \mathcal{H}} + 5s - f_{\text{evo}}^N \right) \left(\Psi + \mathbf{V} \cdot \mathbf{n} + \int_0^{r_S} dr (\Phi' + \Psi') \right) \\ &+ \frac{2 - 5s}{2r_S} \int_0^{r_S} dr \left[2 - \frac{r_S - r}{r} \Delta_\Omega \right] (\Phi + \Psi) \end{aligned}$$

z-dependent angular power spectrum

Multipole expansion

$$\Delta(\mathbf{n}, z) = \sum_{\ell m} a_{\ell m}(z) Y_{\ell m}(\mathbf{n})$$

Power spectrum

$$c_{\ell}(z, z') = \langle a_{\ell m}(z) a_{\ell m}^*(z') \rangle$$

$$c_{\ell}(z_1, z_2) = \langle a_{\ell m}(z_1) a_{\ell m}(z_2) \rangle = 4\pi \int \frac{dk}{k} \mathcal{P}(k) \Delta_{\ell}(z_1, k) \Delta_{\ell}(z_2, k)$$

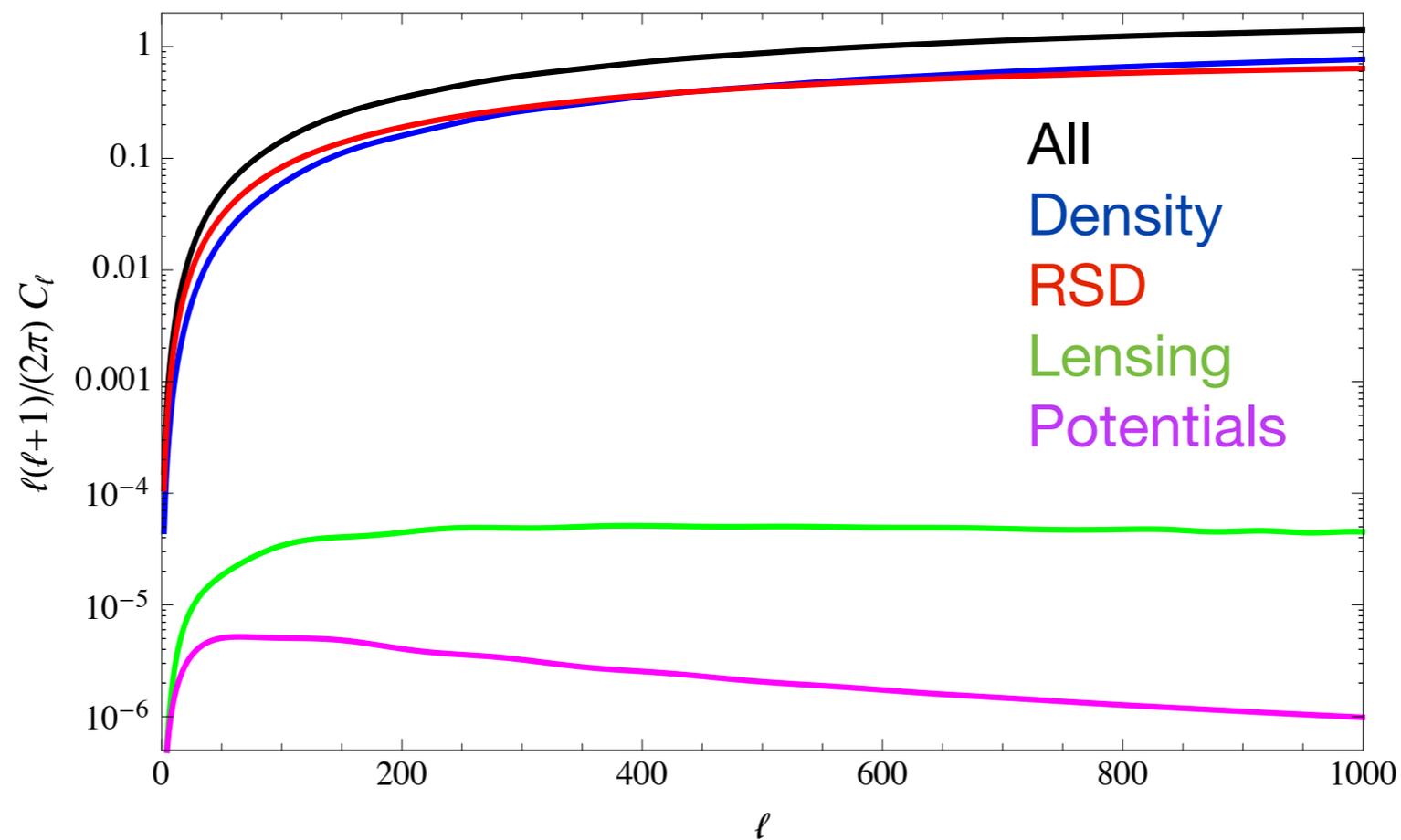
Correlation function

$$\xi(z, z', \theta) = \langle \Delta(\mathbf{n}, z) \Delta(\mathbf{n}', z') \rangle = \sum_{\ell} \frac{2\ell + 1}{4\pi} c_{\ell}(z, z') P_{\ell}(\cos \theta)$$

z-dependent angular power spectrum

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z-dependent angular power spectrum

Window Function

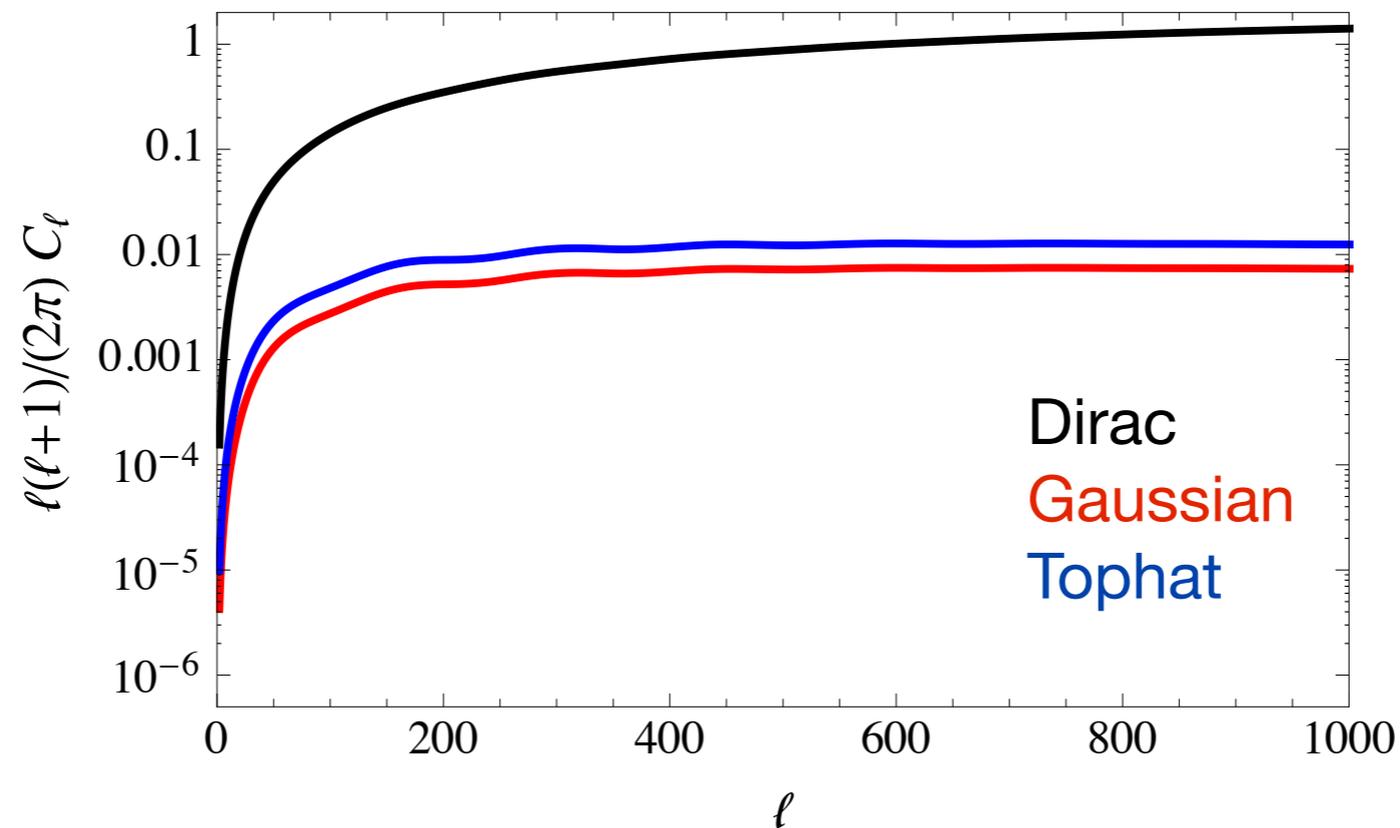
$$\Delta_{\ell}^i(k) = \int dz \frac{dN}{dz} W_i(z) \Delta_{\ell}(z, k) \quad c_{\ell}^{ij} = 4\pi \int \frac{dk}{k} \mathcal{P}(k) \Delta_{\ell}^i(k) \Delta_{\ell}^j(k)$$

$$i = j$$

Auto-Correlation

$$i \neq j$$

Cross-Correlation



z-dependent angular power spectrum

Window Function

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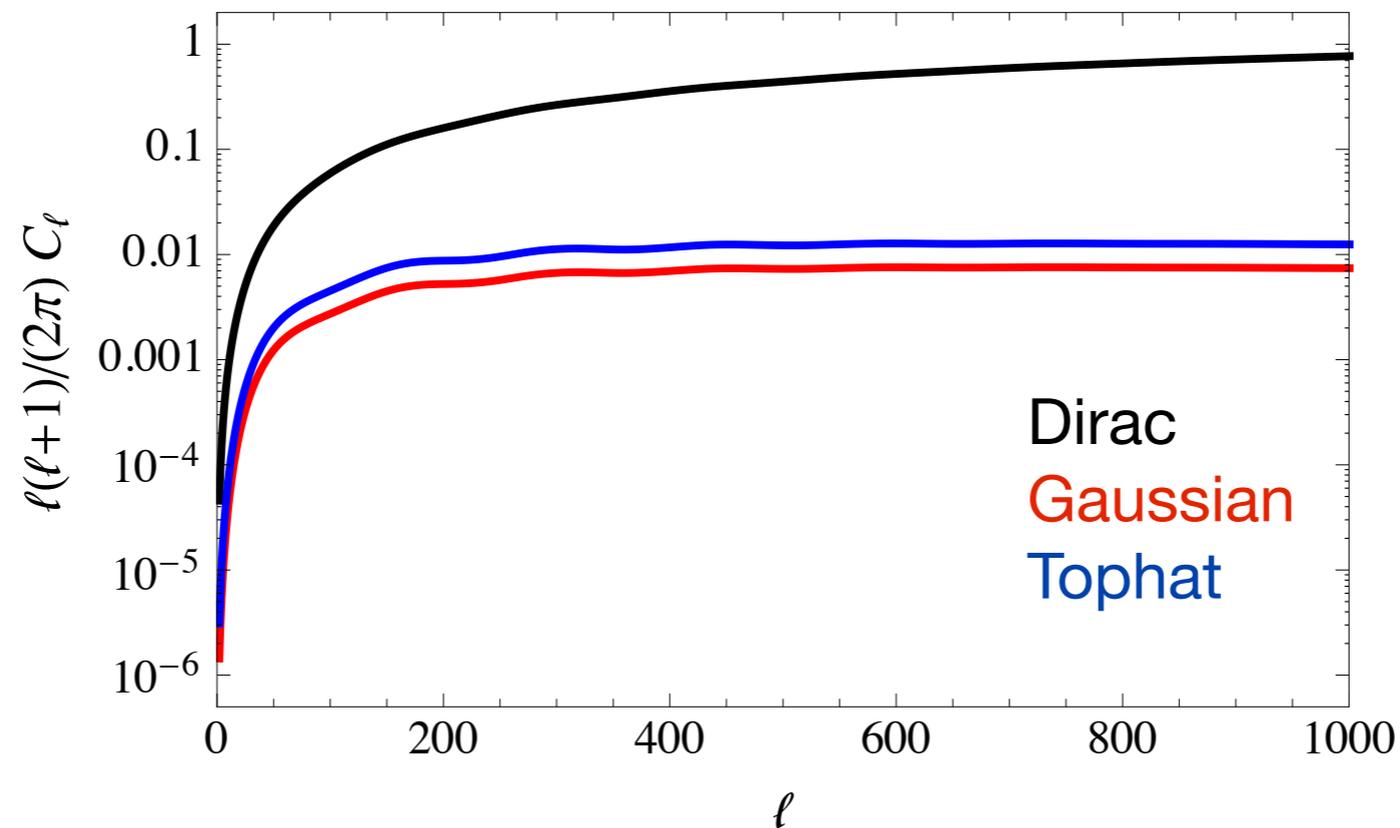
$$i = j$$

Auto-Correlation

$$i \neq j$$

Cross-Correlation

Density



z-dependent angular power spectrum

Window Function

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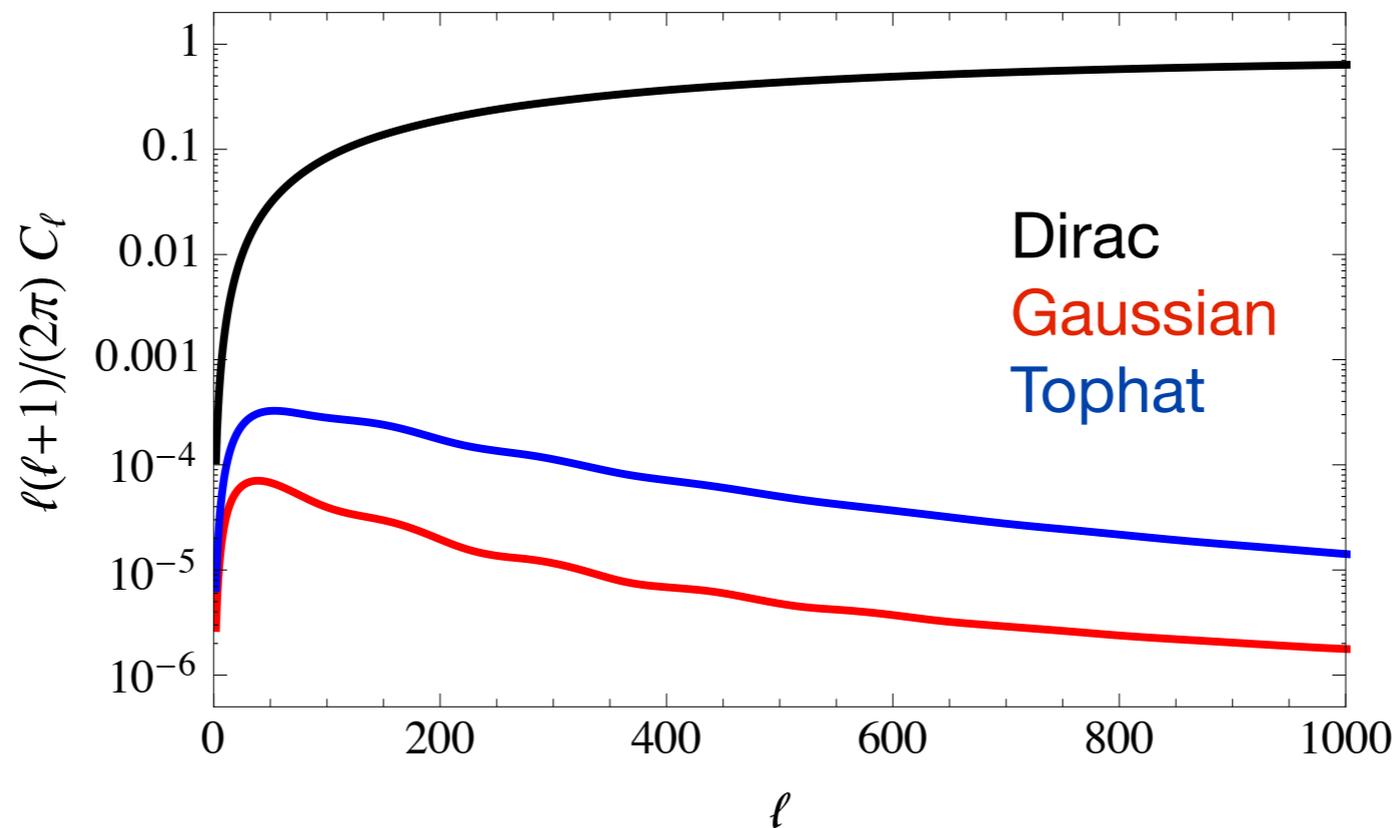
$$i = j$$

Auto-Correlation

$$i \neq j$$

Cross-Correlation

RSD



z-dependent angular power spectrum

Window Function

$$\Delta_{\ell}^i(k) = \int dz \frac{dN}{dz} W_i(z) \Delta_{\ell}(z, k) \quad c_{\ell}^{ij} = 4\pi \int \frac{dk}{k} \mathcal{P}(k) \Delta_{\ell}^i(k) \Delta_{\ell}^j(k)$$

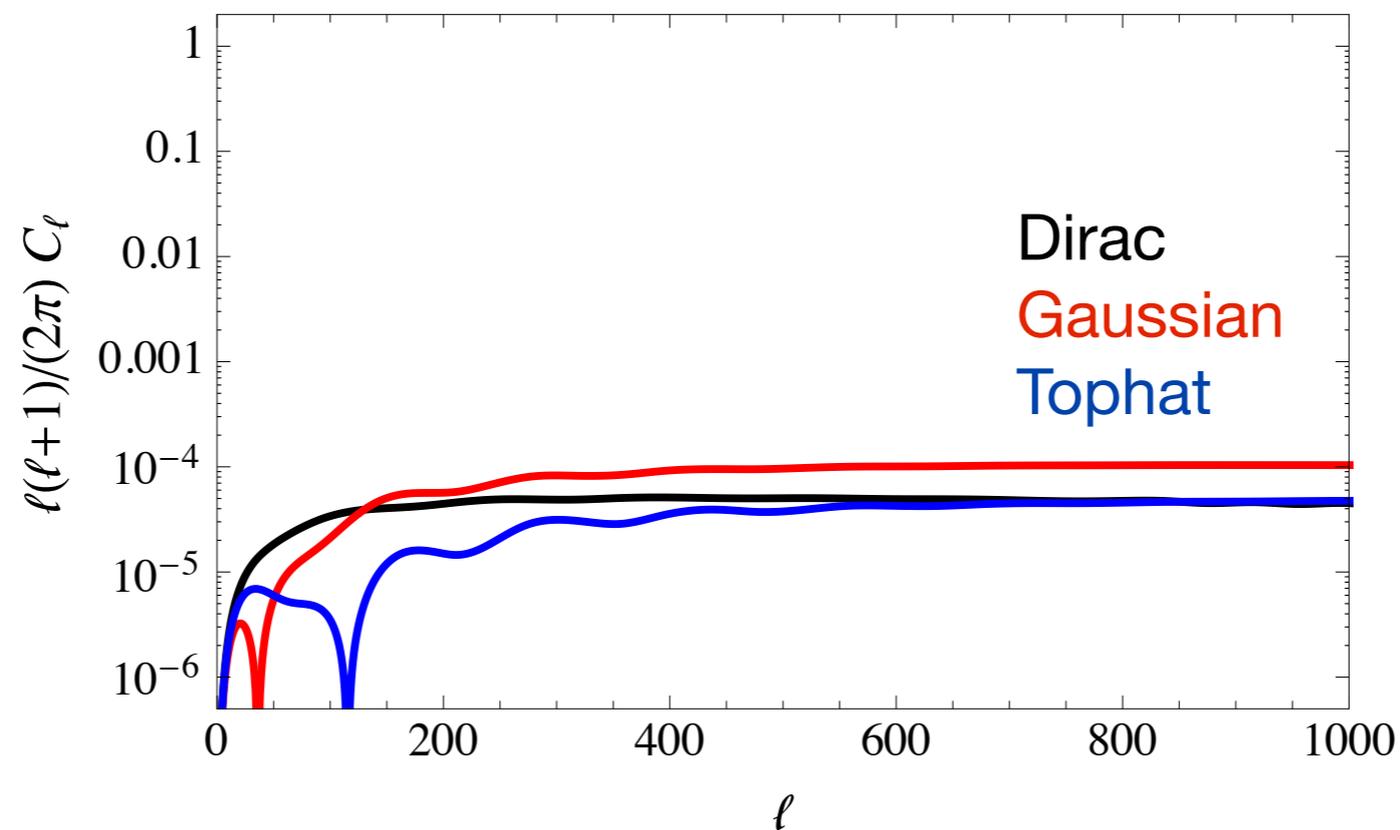
$$i = j$$

Auto-Correlation

$$i \neq j$$

Cross-Correlation

Lensing



z-dependent angular power spectrum

Window Function

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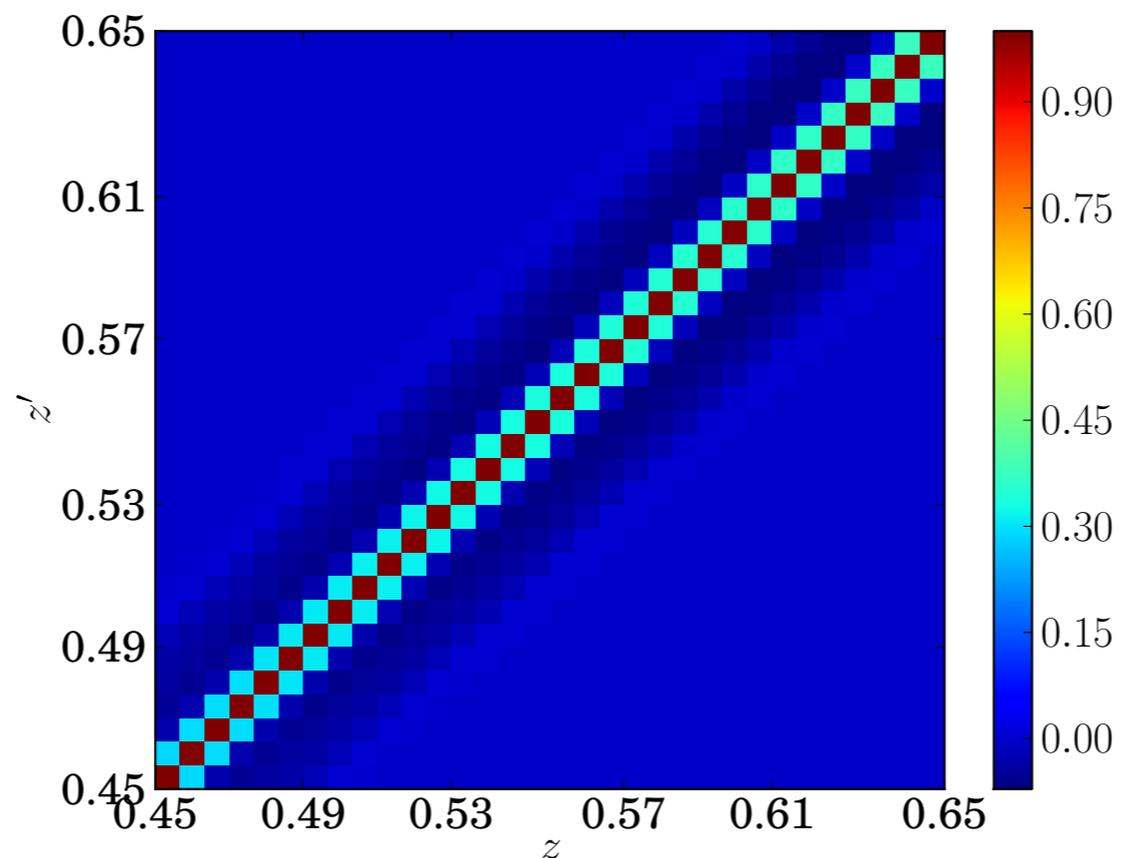
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Cross-Correlation

$$\ell = 71$$



Fisher Matrix Analysis

How good can a future survey determine the cosmological parameters?

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$$F_{\alpha\beta} = \sum_{\ell, (ij), (pq)} \frac{\partial c_{\ell}^{ij}}{\partial \lambda_{\alpha}} \frac{\partial c_{\ell}^{pq}}{\partial \lambda_{\beta}} \text{Cov}_{\ell, (ij), (pq)}^{-1}$$

(assuming gaussian likelihood)

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Cosmological parameters λ_{α}

Covariance matrix

$$\text{Cov}_{[\ell, \ell'][(ij), (pq)]} = \delta_{\ell, \ell'} \frac{C_{\ell}^{\text{obs}, ip} C_{\ell}^{\text{obs}, jq} + C_{\ell}^{\text{obs}, iq} C_{\ell}^{\text{obs}, jp}}{f_{\text{sky}} (2\ell + 1)}$$

Asorey, Crocce, Gaztañaga, Lewis [arXiv:1207.6487]

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Observed spectrum

$$C_{\ell}^{\text{obs}, ij} = C_{\ell}^{ij} + \frac{\delta_{ij}}{n(i)}$$

$n(i)$

number of galaxies per steradian in the i -th redshift bin

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Observed spectrum

$$C_{\ell}^{\text{obs}, ij} = C_{\ell}^{ij} - \frac{\delta_{ij}}{n(i)}$$

$n(i)$

number of galaxies per steradian in the i -th redshift bin

Shot-noise

Fisher Matrix Analysis

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Figure of Merit

$$\text{FoM} = [\det(F^{-1})]^{-1/2}$$

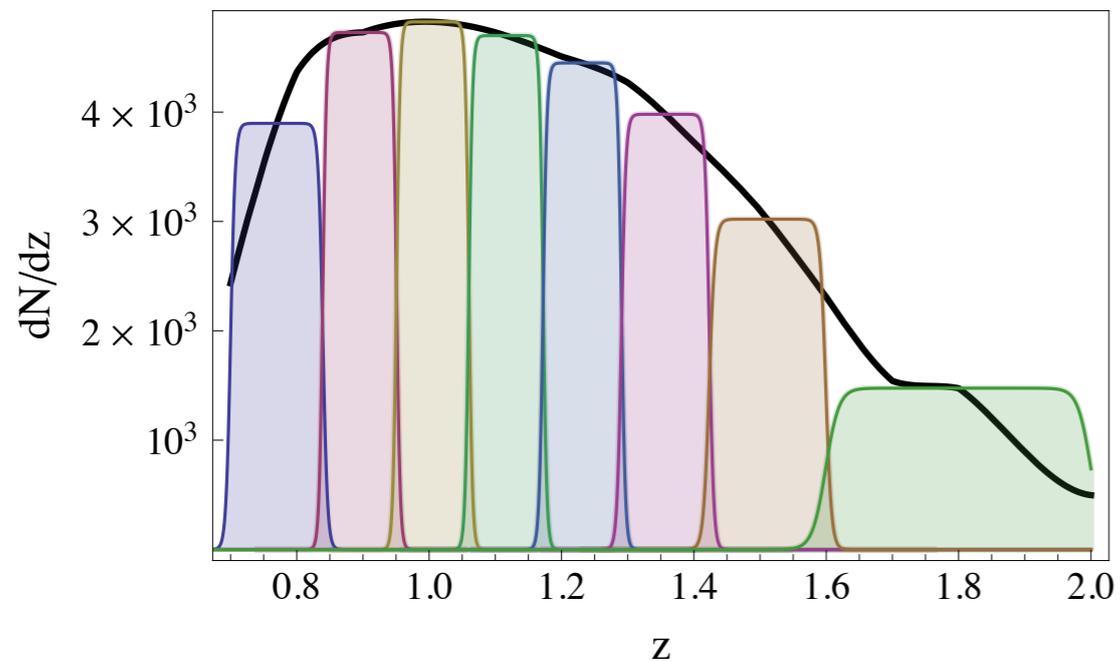
$$\text{FoM}_{\text{fixed}} = [\det(\hat{F}^{-1})]^{-1/2}$$

$$\text{FoM}_{\text{marginalized}} = [\det(\widehat{F}^{-1})]^{-1/2}$$

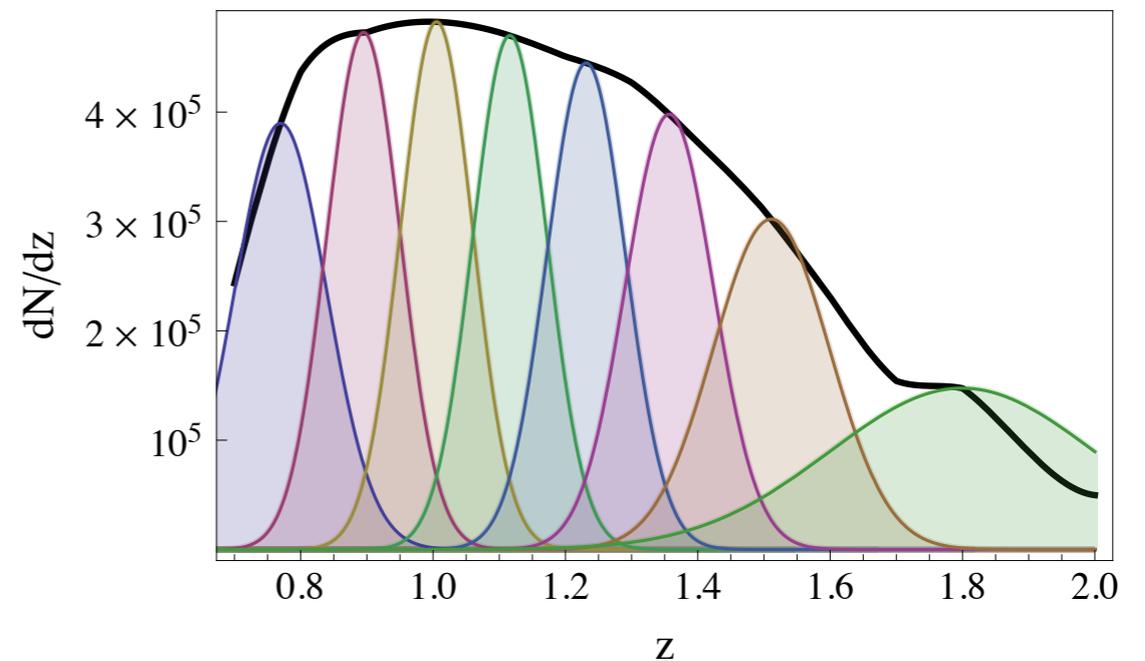
Binning Strategy

To recover the 3D information we need to split the redshift range in many bins and to consider the cross-correlations between different redshift bins.

Spectroscopic Survey



Photometric Survey



The redshift bins are chosen such that there is the same number of galaxy for each bin.

Non-linearity scales

We can not correlate two galaxies separated by a distance shorter than the non-linear scale λ_{\min}

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Longitudinal distance $\frac{z_i - z_j}{H(\bar{z})}$

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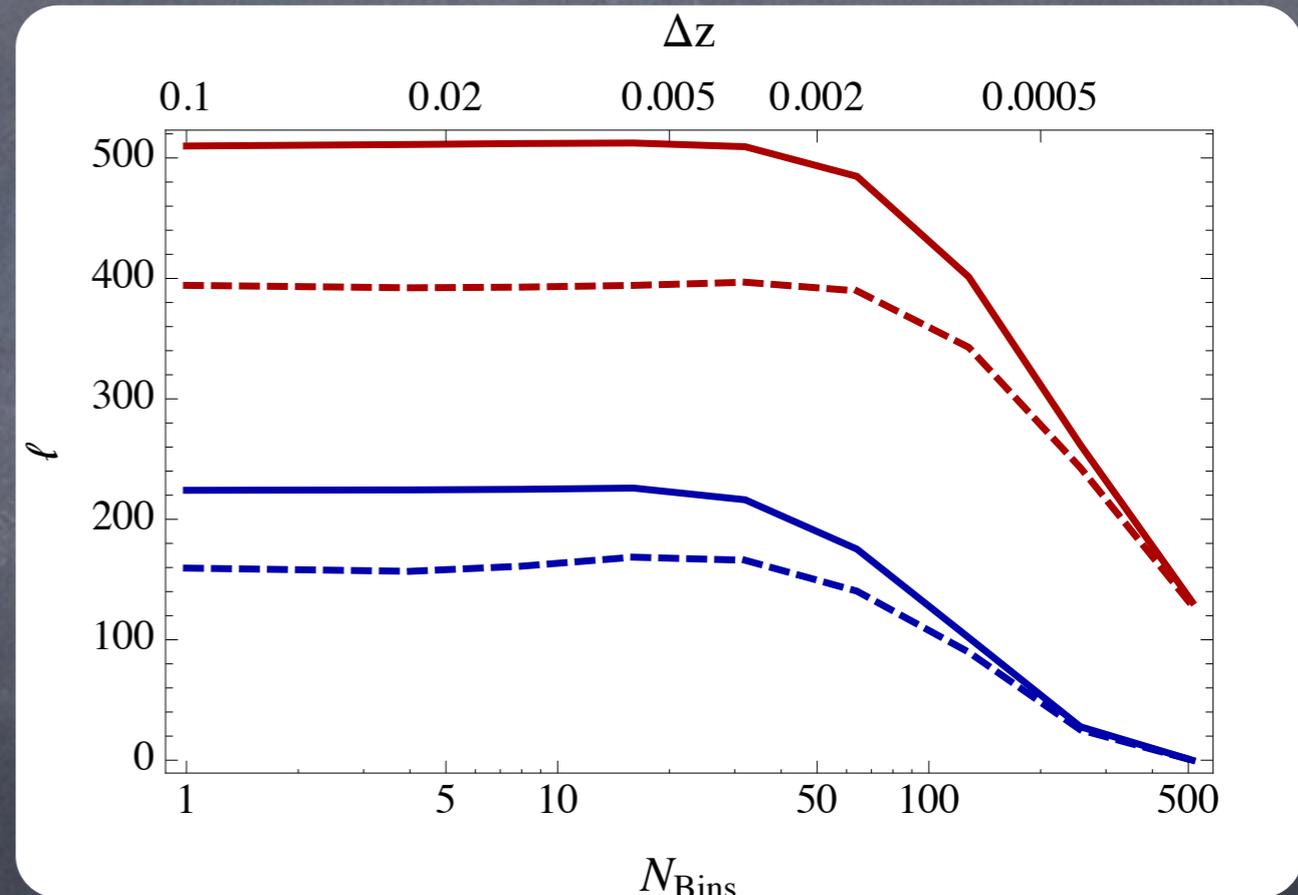
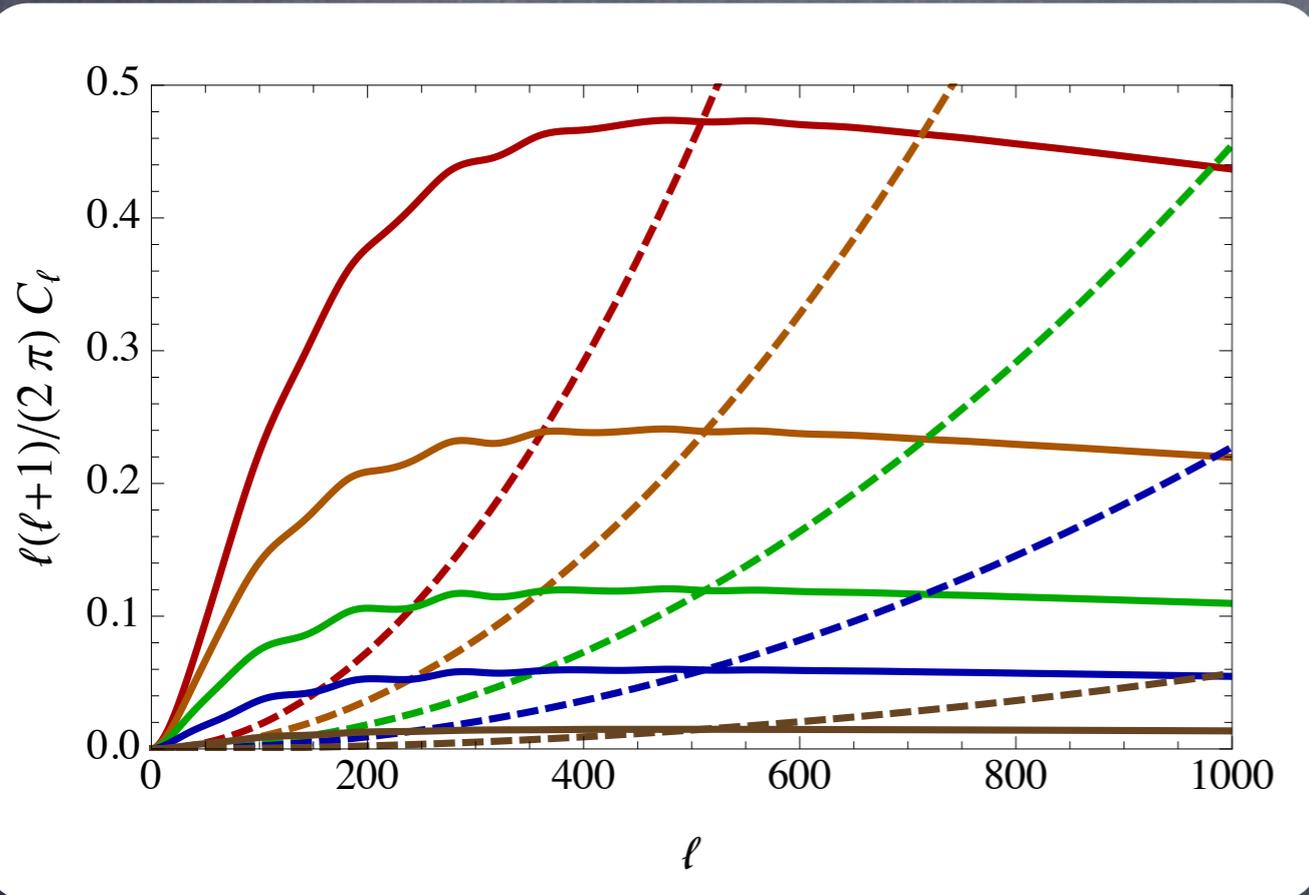
Longitudinal distance $\frac{z_i - z_j}{H(\bar{z})}$

$$\ell_{\max}^{ij} = \begin{cases} 2\pi d(\bar{z})/\lambda_{\min} & \text{if } \delta z_{ij} \leq 0, \\ \frac{2\pi d(\bar{z})}{\sqrt{\lambda_{\min}^2 - \left(\frac{\delta z_{ij}}{H(\bar{z})}\right)^2}} & \text{if } 0 < \frac{\delta z_{ij}}{H(\bar{z})} < \lambda_{\min}, \\ \infty & \text{otherwise.} \end{cases}$$

Cosmological Parameter Forecast

Increasing the number of bins, we get a better redshift resolution.

But shot-noise increases too.



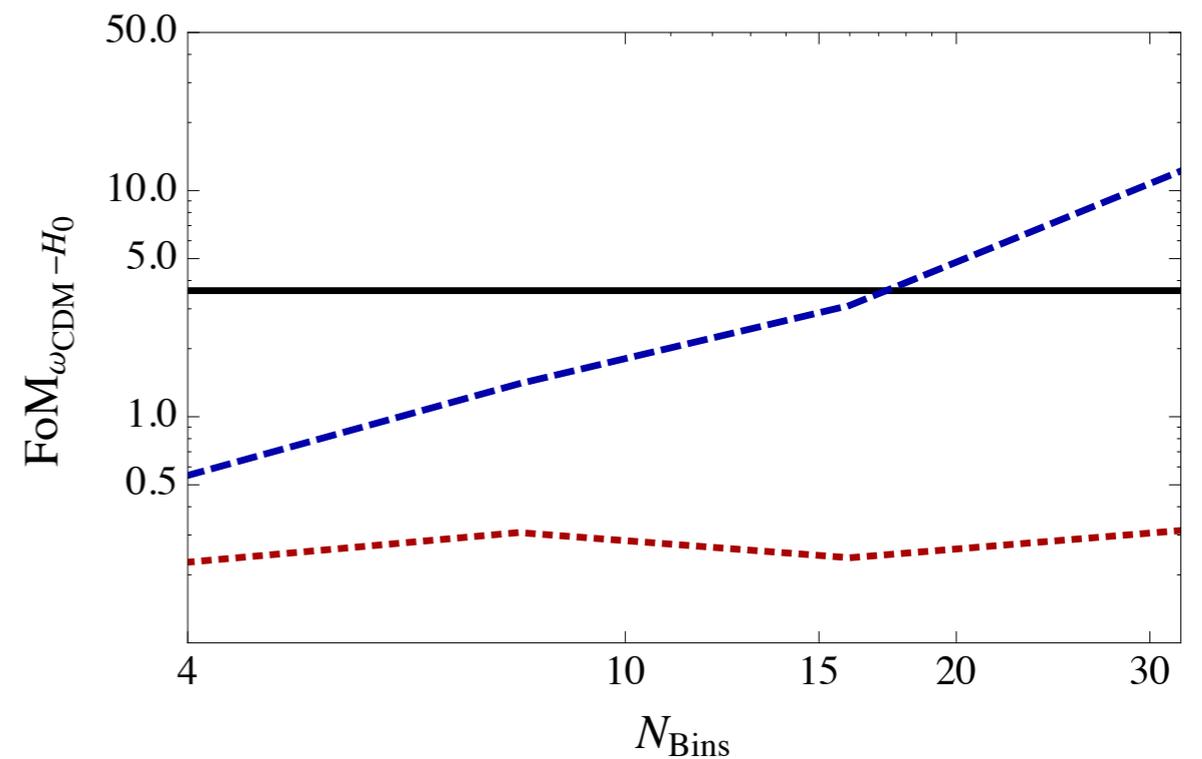
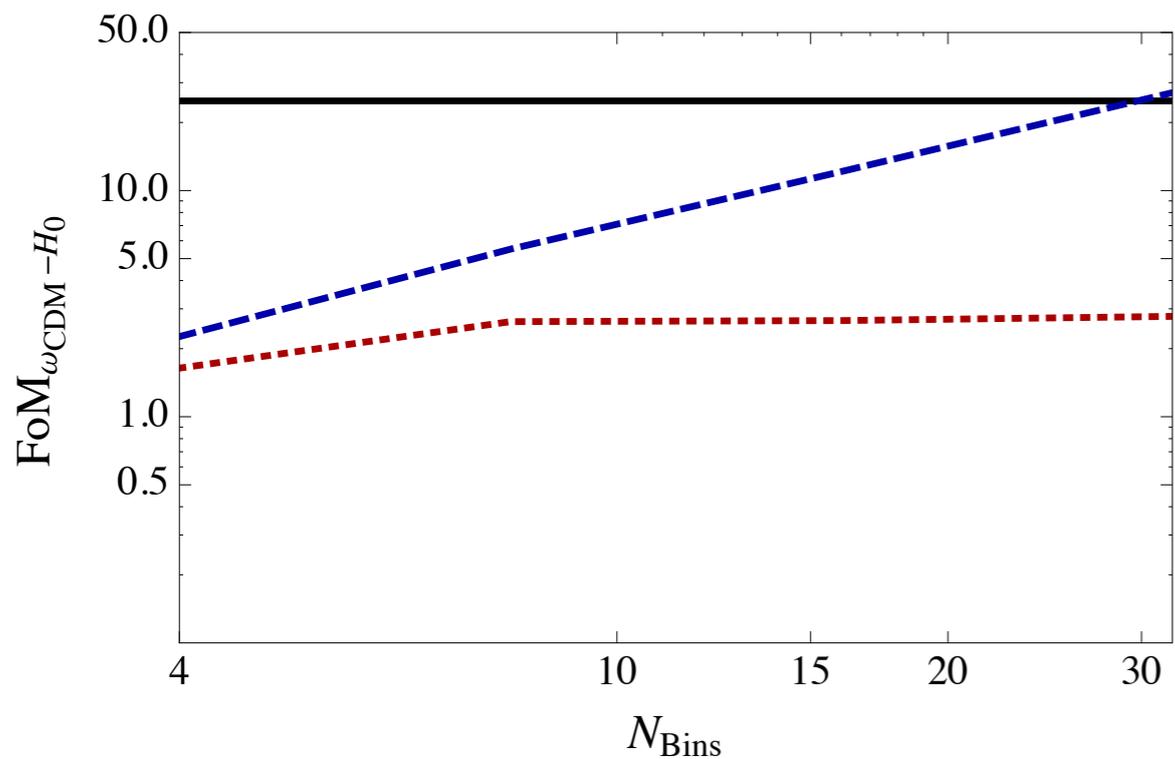
Since also the signal increases, shot-noise starts dominating for smaller bin size than naively expected.

Cosmological Parameter Forecast

spectroscopic DES-like

$$\lambda_{\min} = 34 \text{ Mpc}/h$$

$$\lambda_{\min} = 68 \text{ Mpc}/h$$



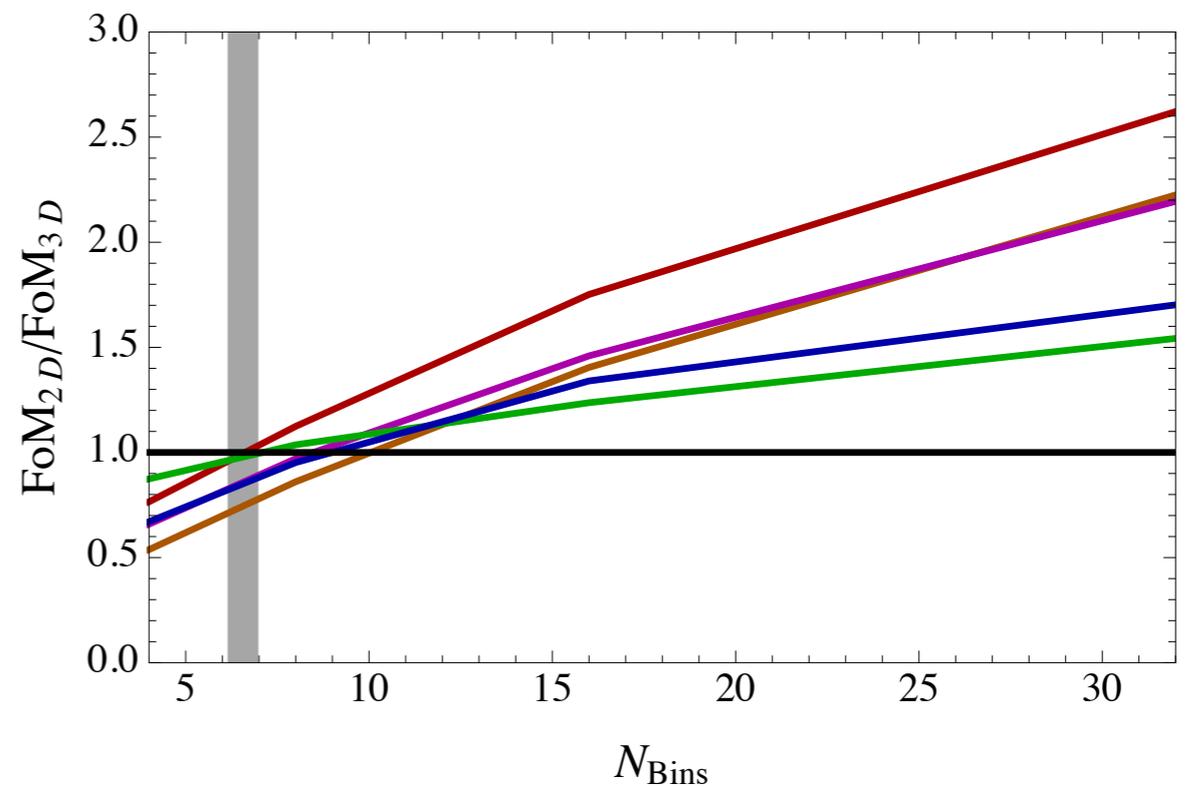
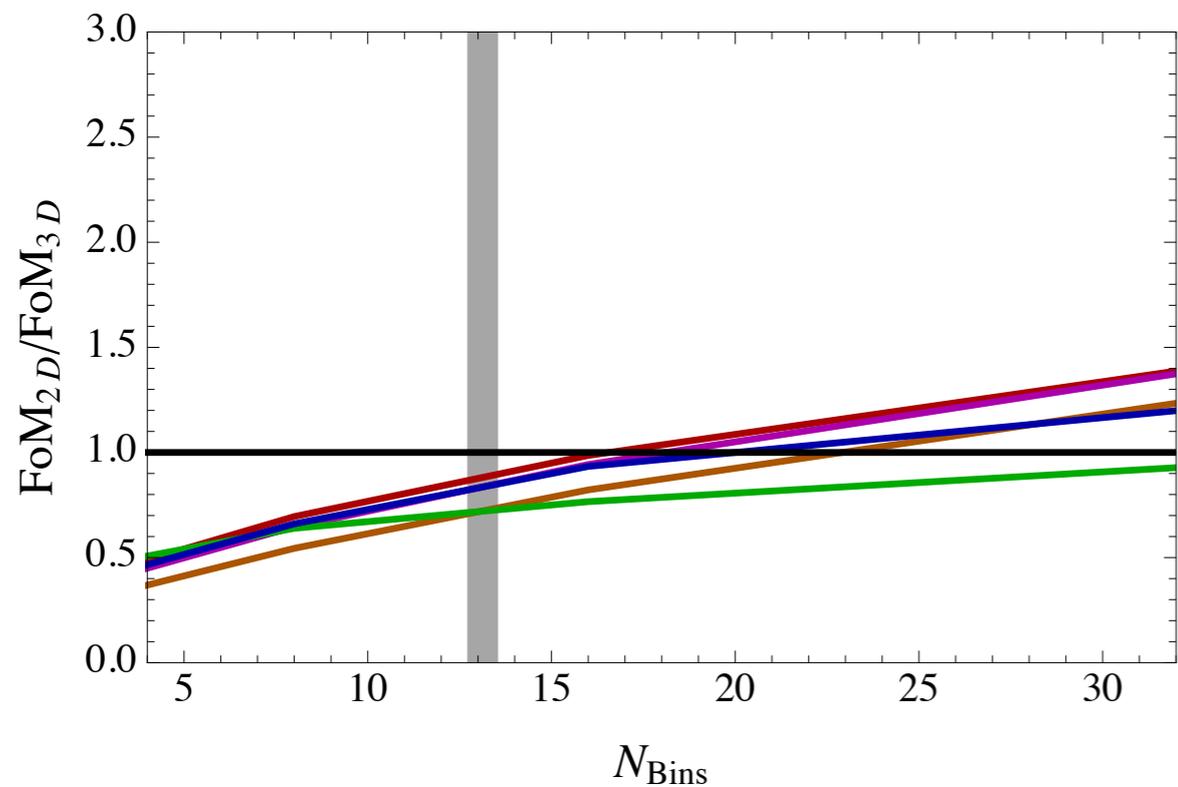
only redshift bins auto-correlations
with redshift bins cross-correlations

Cosmological Parameter Forecast

2D vs 3D

$$\lambda_{\min} = 34 \text{ Mpc}/h$$

$$\lambda_{\min} = 68 \text{ Mpc}/h$$

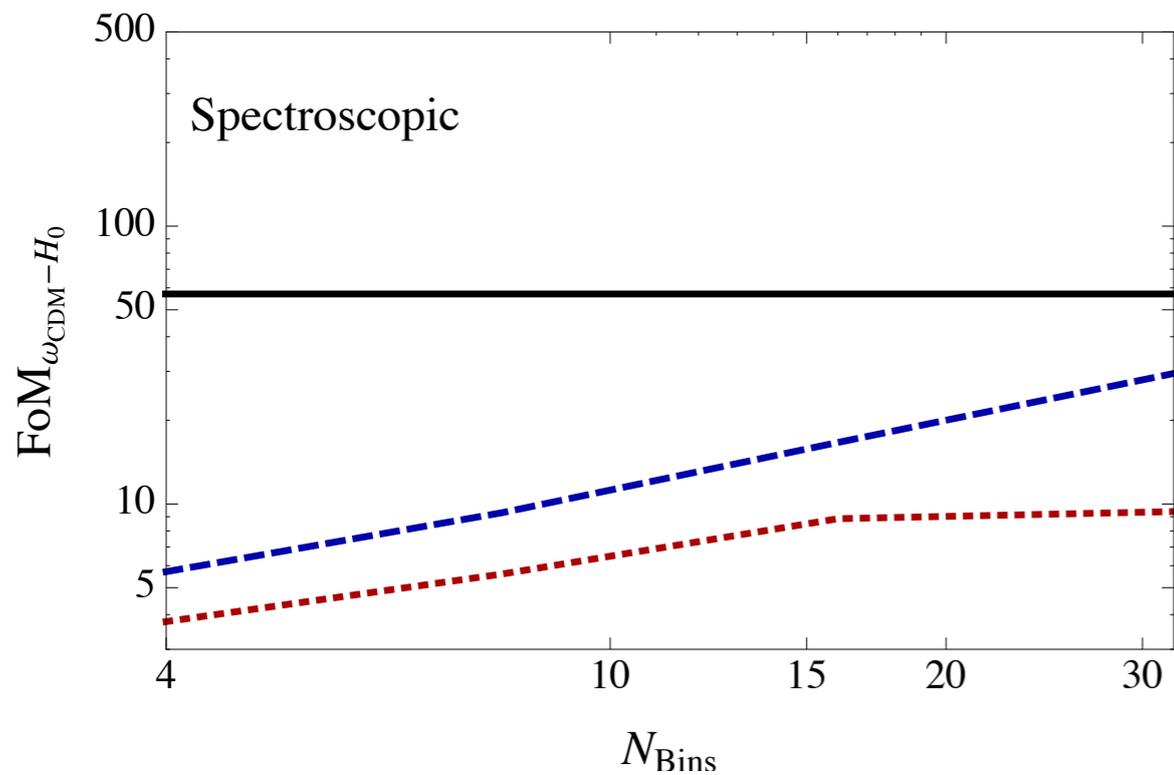


ω_b ω_{CDM} n_s H_0 A_s

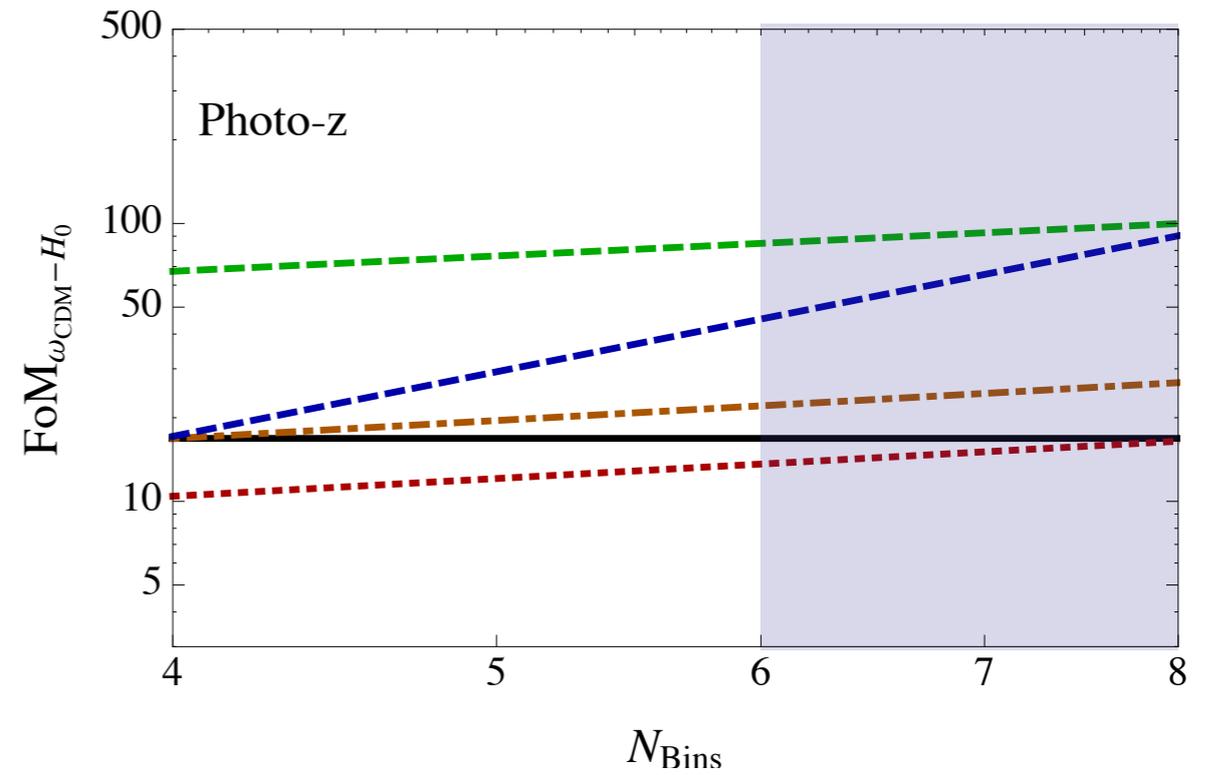
Cosmological Parameter Forecast

Euclid

Spectroscopic Survey



Photometric Survey

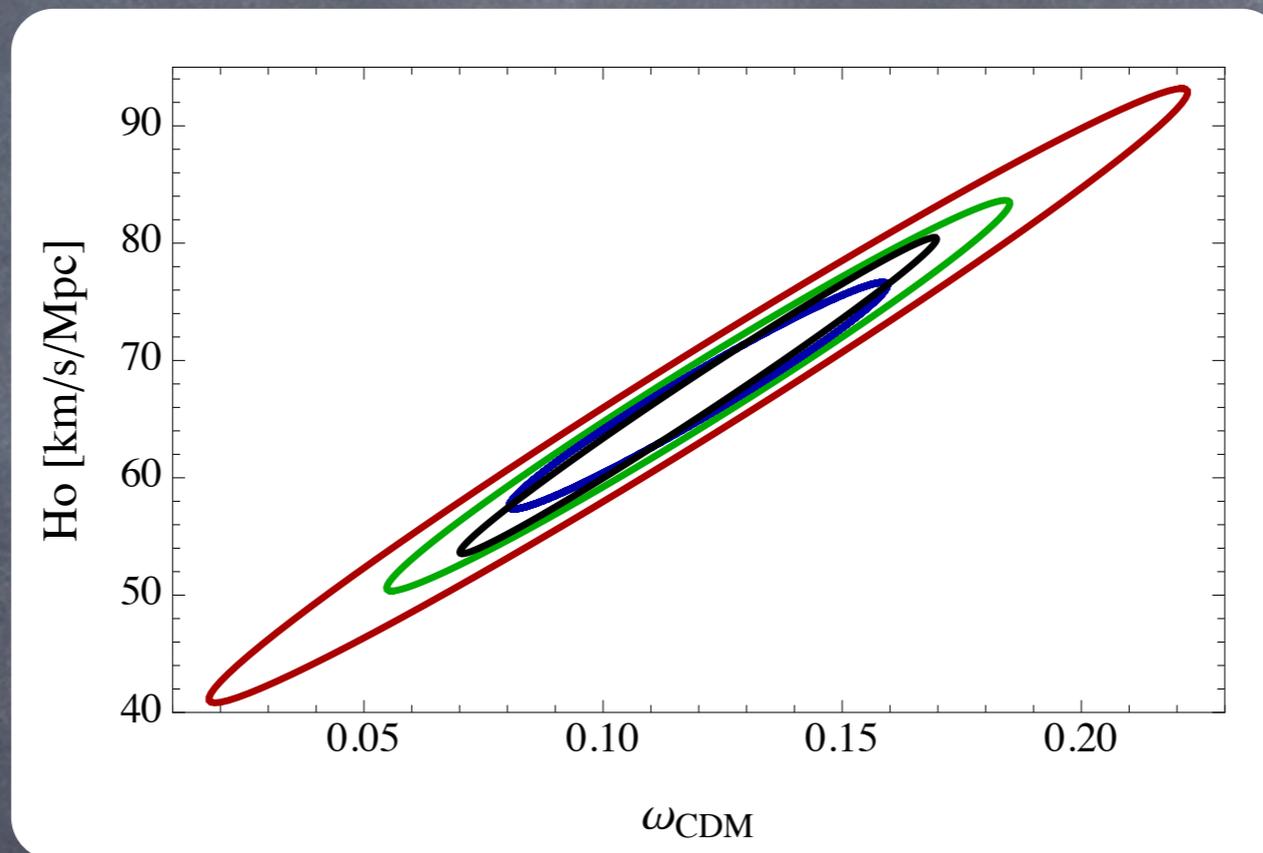


$$\lambda_{\min} = 68 \text{ Mpc}/h$$

only redshift bins auto-correlations
with redshift bins cross-correlations

Cosmological Parameter Forecast

spectroscopic DES-like

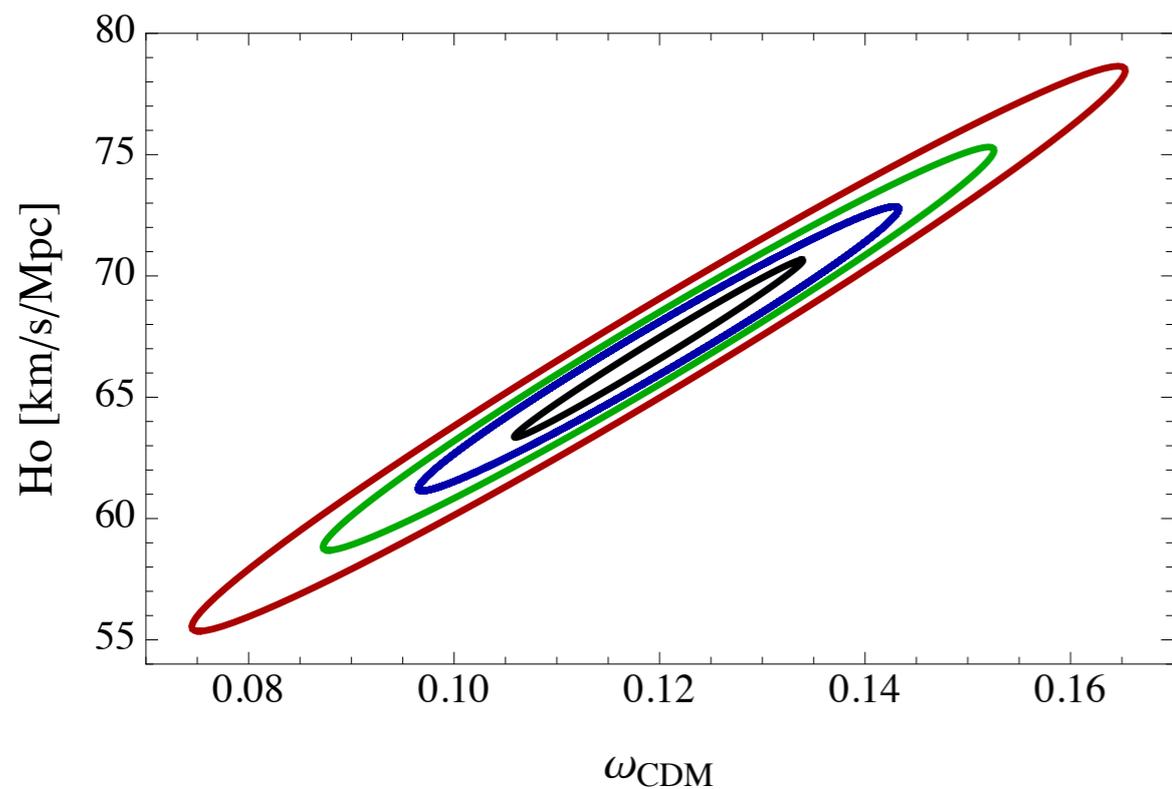


4 bins, 8 bins, 16 bins, 32 bins

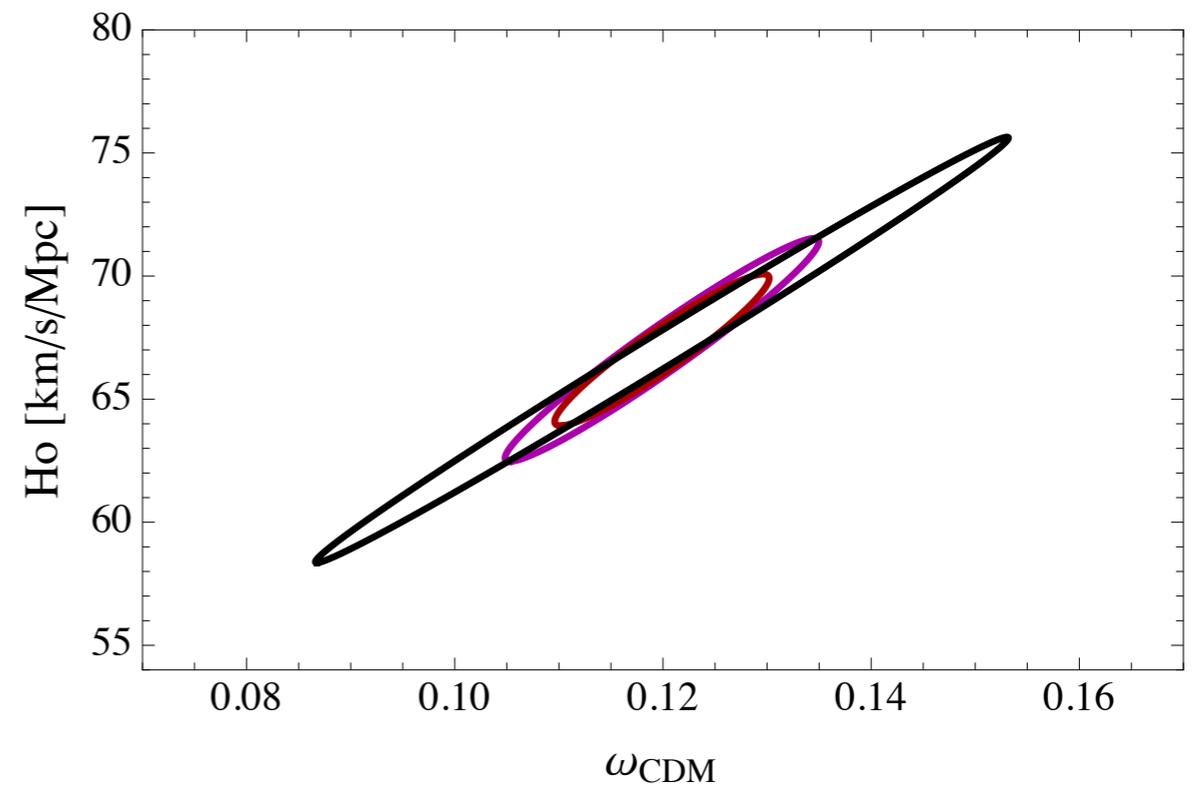
Cosmological Parameter Forecast

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Spectroscopic Survey



Photometric Survey



4 bins, 8 bins, 16 bins, 32 bins

Signal to Noise for different effects

Different terms encode
different information

Density

Redshift-space distortion

Lensing

Potential

Signal to Noise for different effects

Different terms encode
different information

Large scale matter
distribution

Density

Redshift-space distortion

Lensing

Potential

Signal to Noise for different effects

Different terms encode
different information

Cosmic velocity field

Density

Redshift-space distortion

Lensing

Potential

Signal to Noise for different effects

Different terms encode
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Geometry of the universe

Density

Redshift-space distortion

Lensing

Potential

Signal to Noise for different effects

Different terms encode
different information

Density

Redshift-space distortion

Lensing

Potential

Signal to Noise

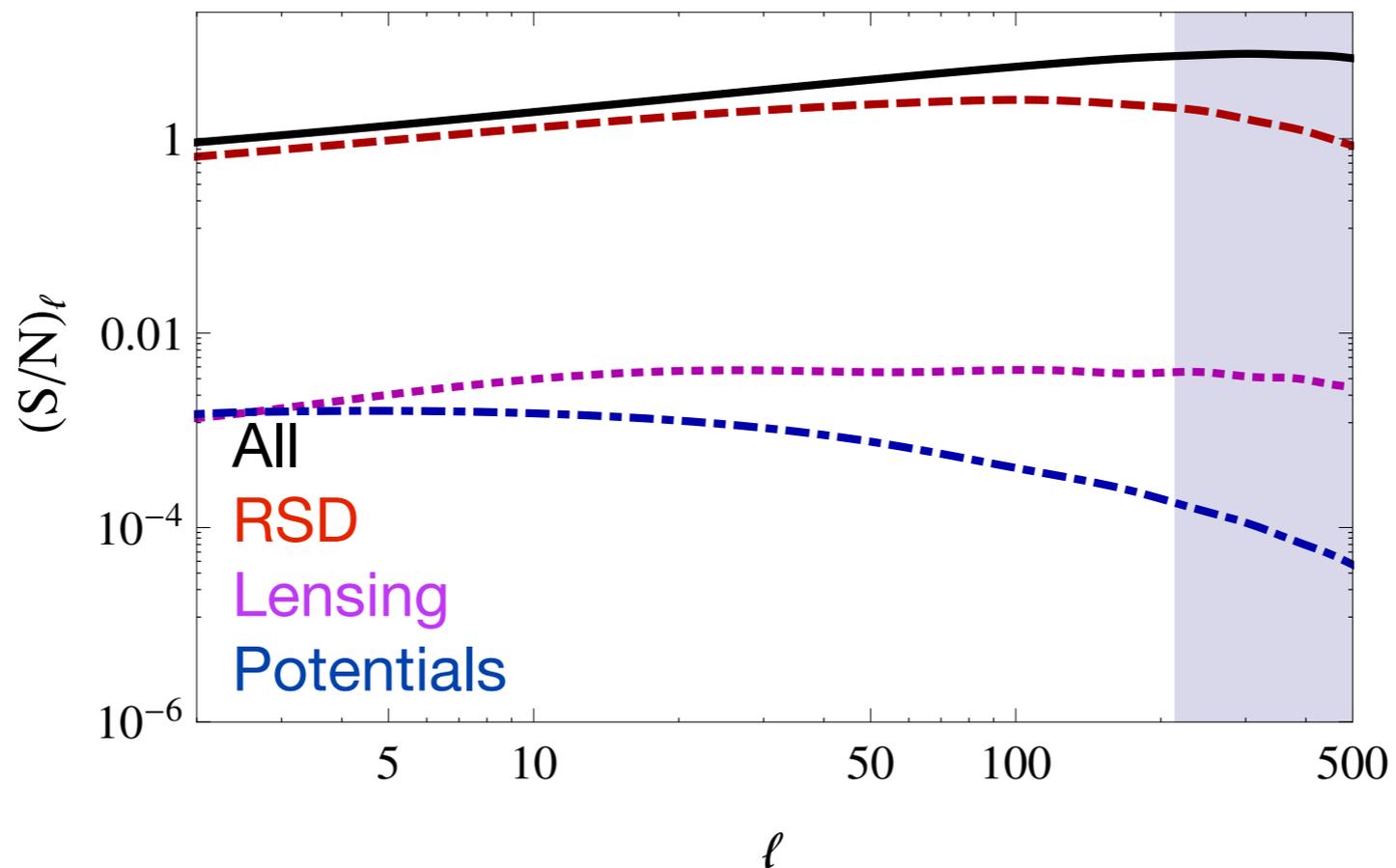
$$\left(\frac{S}{N}\right)_\ell = \frac{C_\ell - \tilde{C}_\ell}{\sigma_\ell}$$
$$\sigma_\ell = \sqrt{\frac{2}{(2\ell + 1) f_{\text{sky}}}} \left(C_\ell + \frac{1}{n}\right)$$

Lensing Potential

Signal to Noise $\left(\frac{S}{N}\right)_\ell = \frac{C_\ell - \tilde{C}_\ell}{\sigma_\ell}$

$$\sigma_\ell = \sqrt{\frac{2}{(2\ell + 1) f_{\text{sky}}}} \left(C_\ell + \frac{1}{n}\right)$$

$\bar{z}=1, \Delta z=0.01$



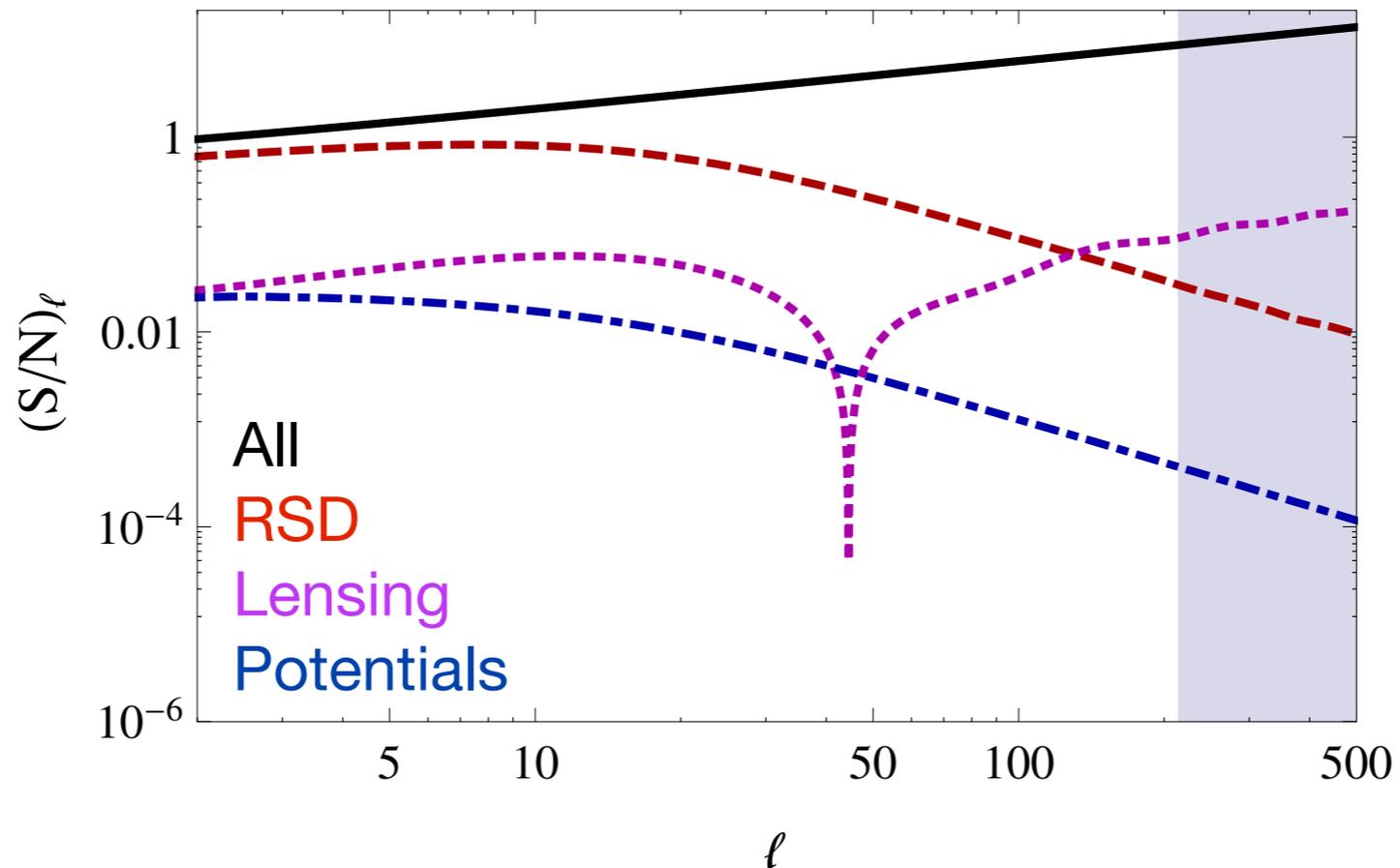
Euclid

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$\bar{z}=1, \Delta z=0.1$



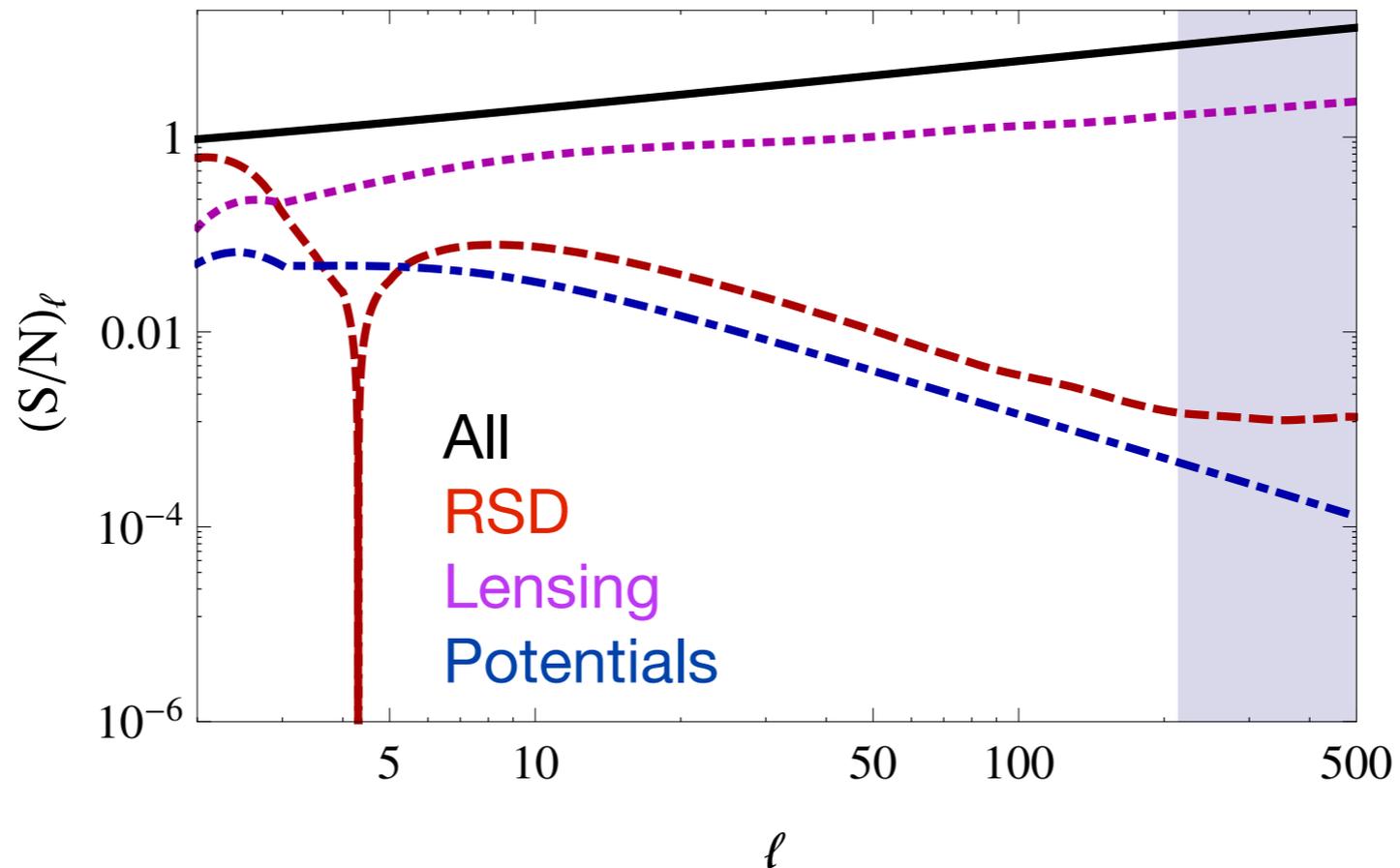
Euclid

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$\bar{z}=1, \Delta z=0.5$

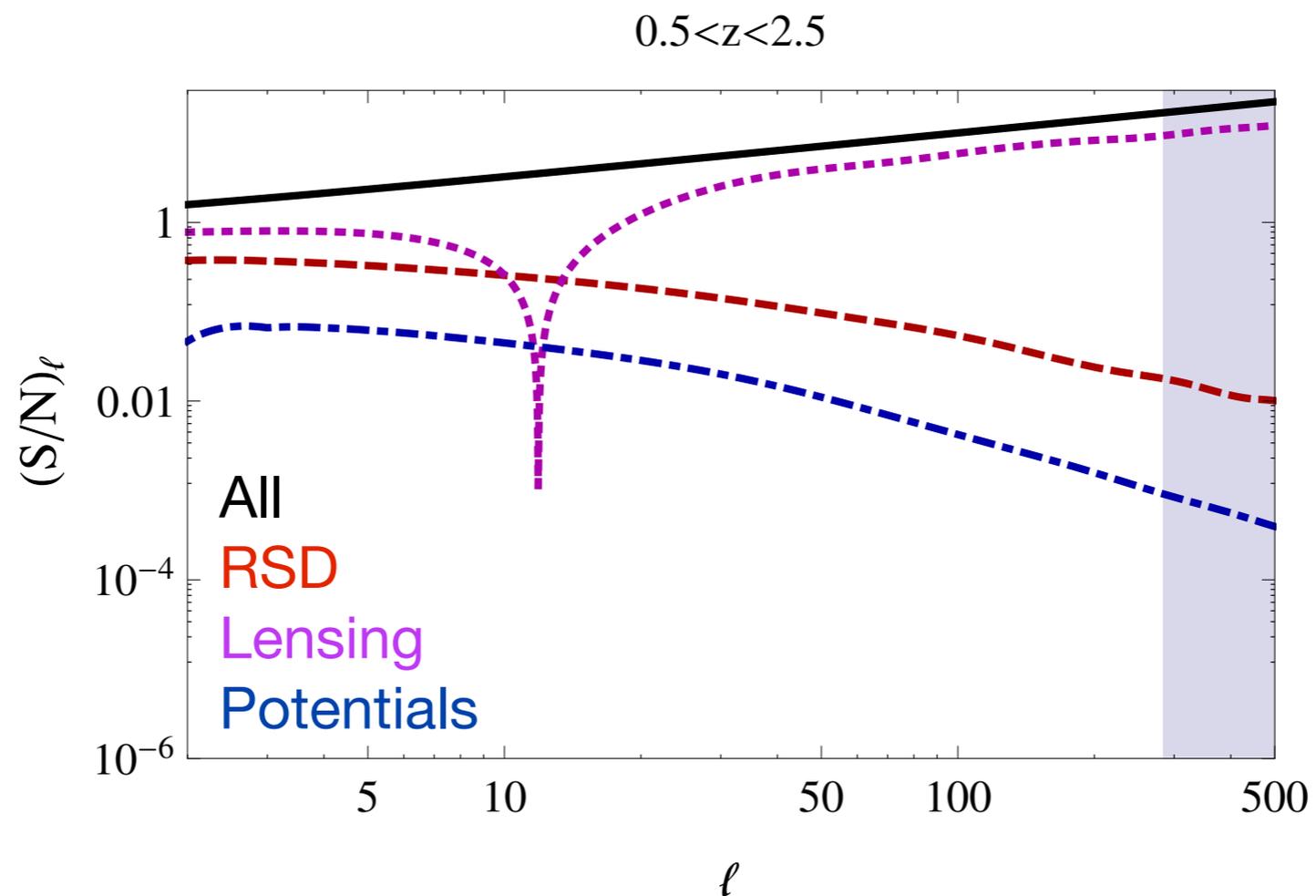


Euclid

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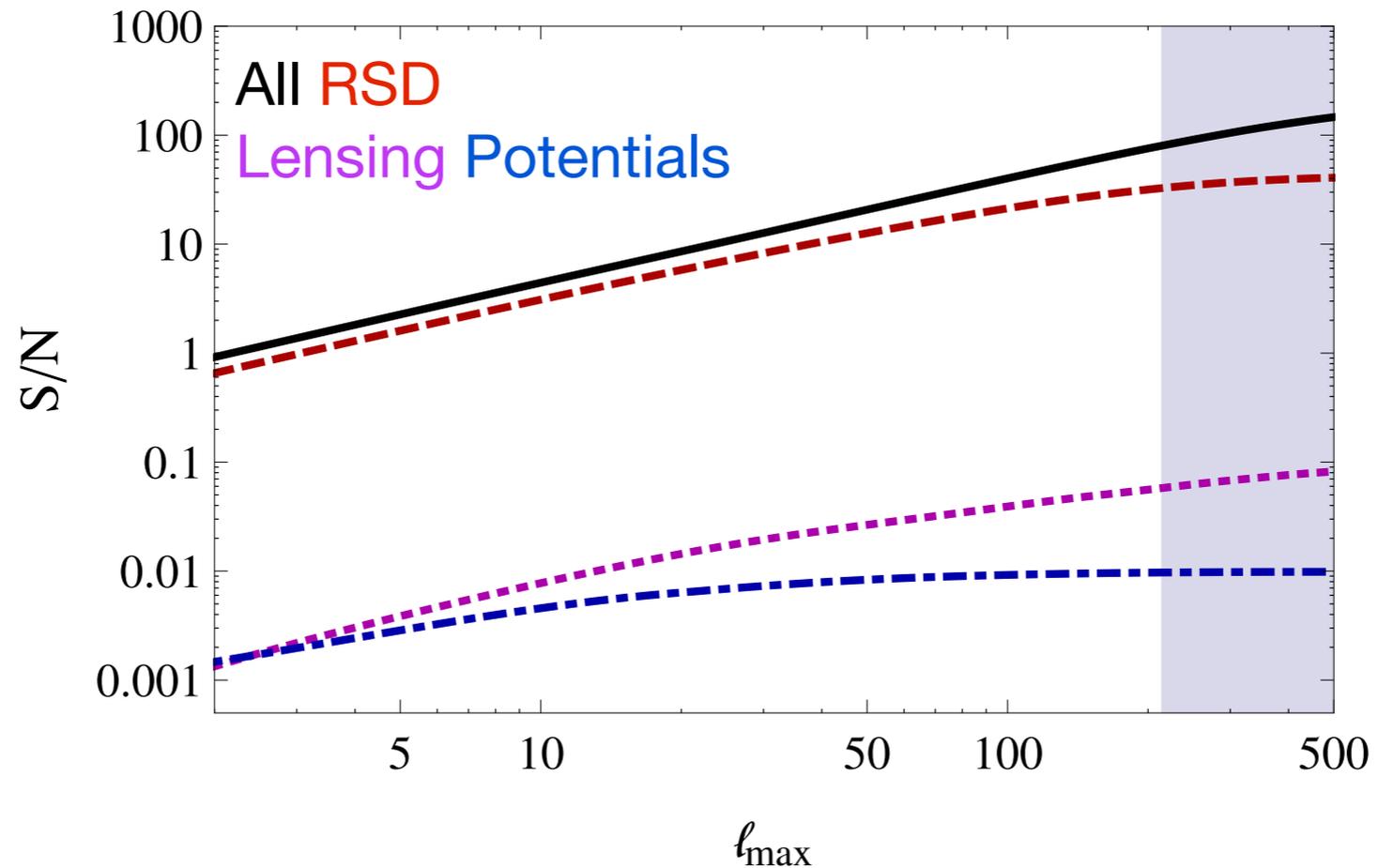


Lensing Potential

Cumulative Signal to Noise

$$\left(\frac{S}{N}\right)^2 = \sum_{\ell=2}^{\ell_{\max}} \left(\frac{C_{\ell} - \tilde{C}_{\ell}}{\sigma_{\ell}}\right)^2$$

$\bar{z}=1, \Delta z=0.01$



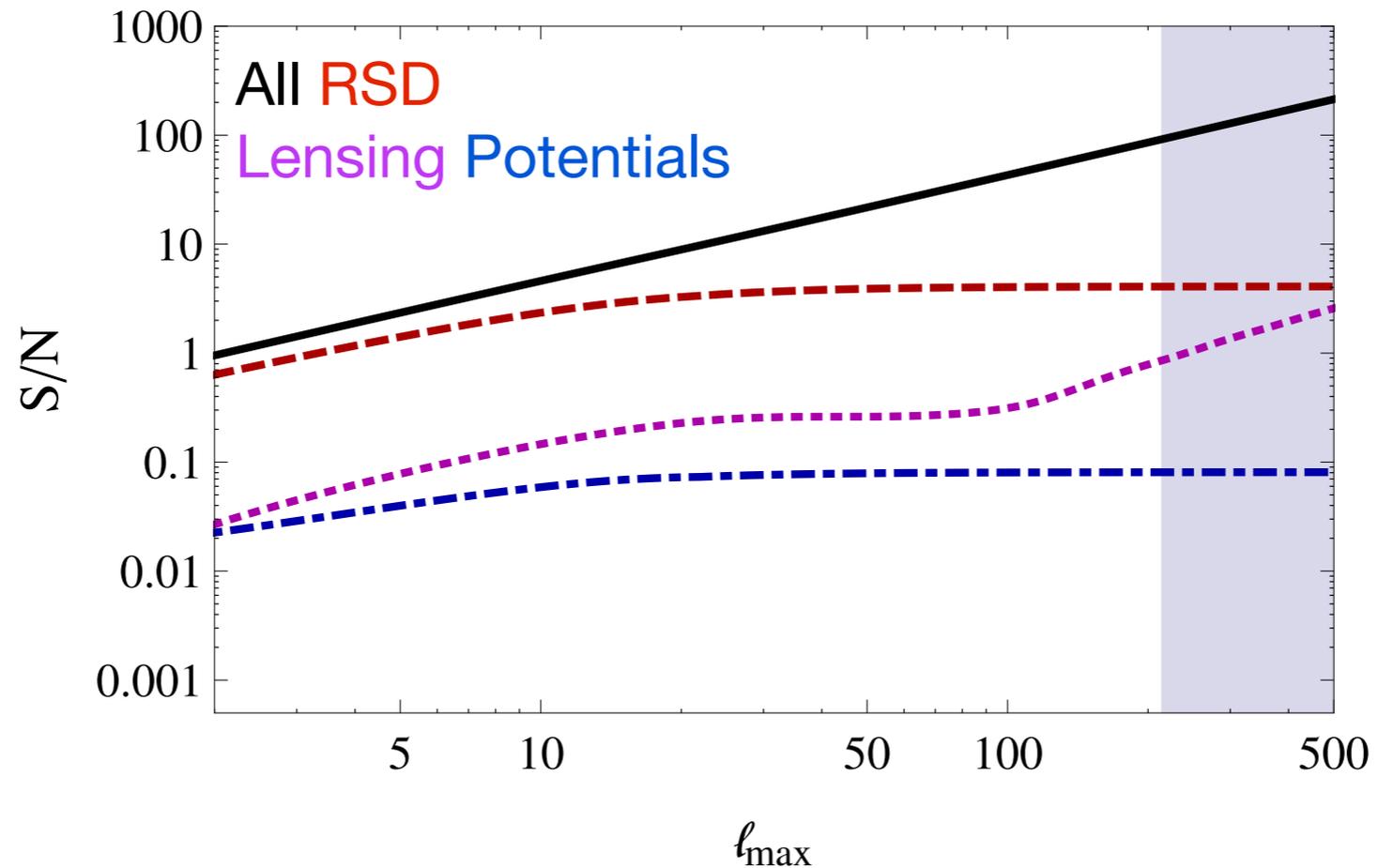
Euclid

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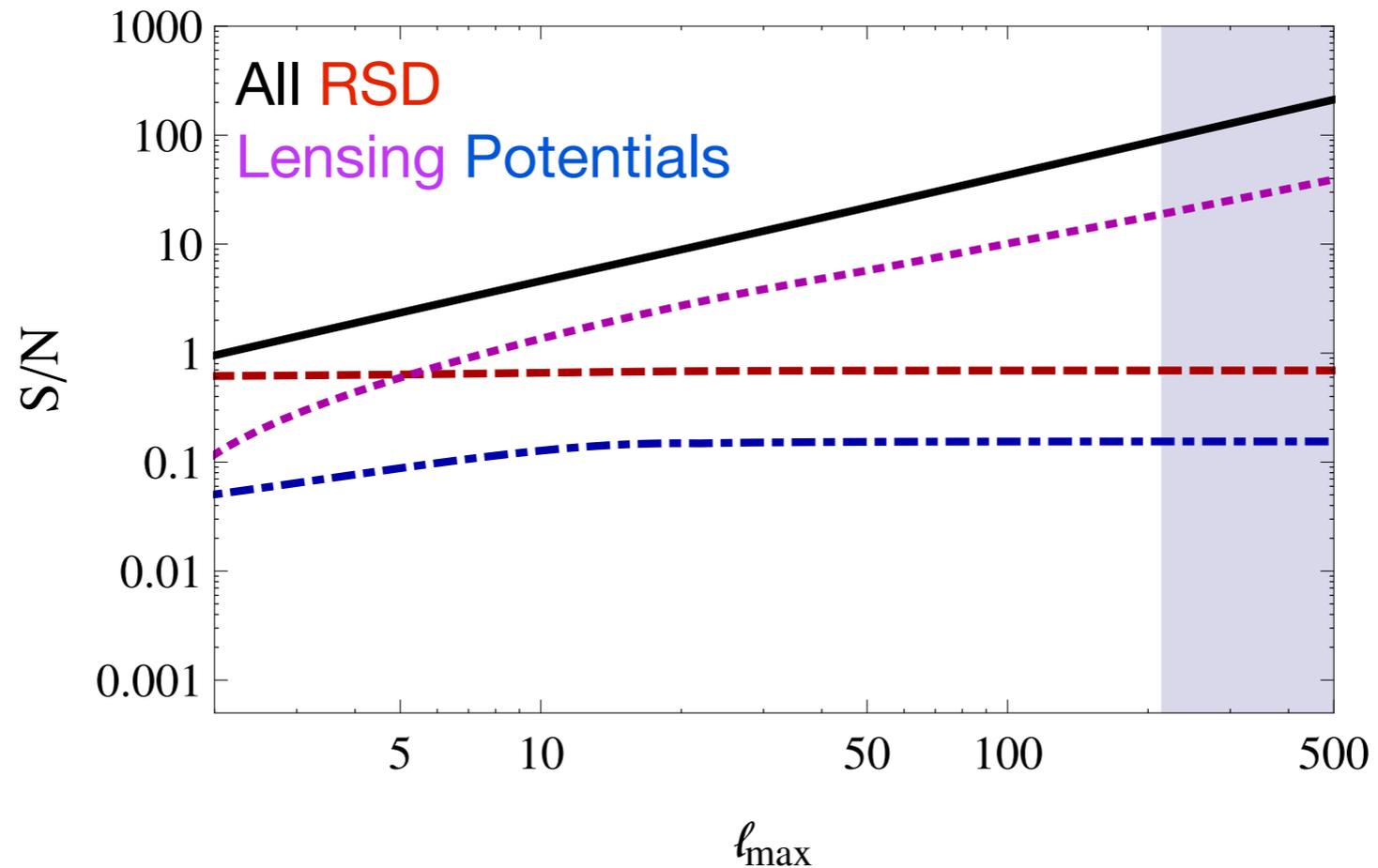
Euclid

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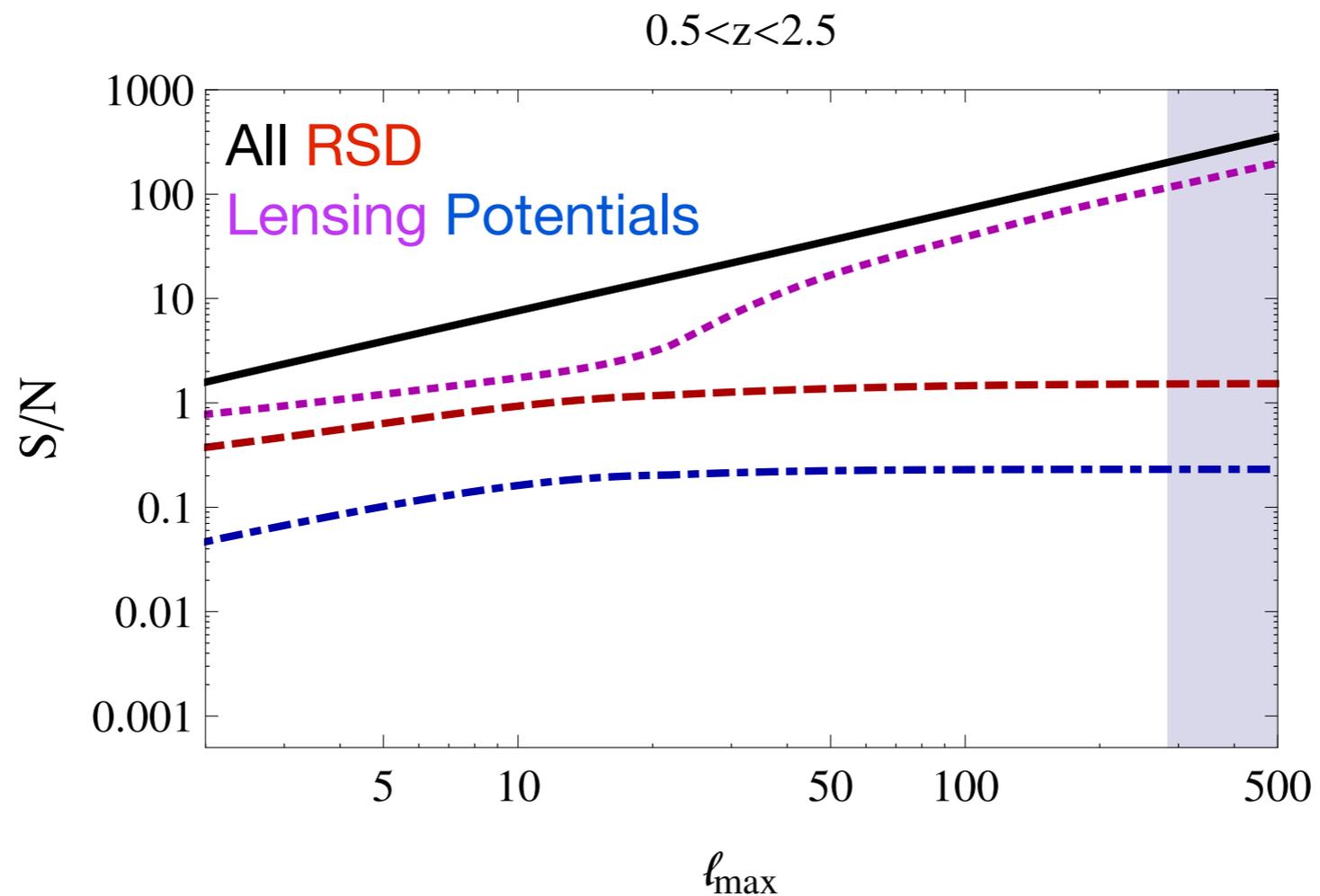


Euclid

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CLASSgal

Modified version of CLASS code to compute the redshift dependent angular power spectra $c_\ell(z_1, z_2)$

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All the relativistic effects:

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- Lensing
- Potential (GR) terms

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MontePython

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CLASSgal

Next..

Main parts of CLASSgal will be merged with CLASS (v2.1):

- cross-correlation with CMB
- implementation for non-flat universe

<http://cosmology.unige.ch/tools>

For CLASS

<http://class-code.net>

Conclusion

So far galaxy surveys have mainly determined $P(k)$.
Easy to measure it, but requires a fiducial model.

Future surveys (like Euclid) will be able to measure the spectra $C_\ell(z, z')$ directly from the data.

These spectra contain information about the matter distribution (density), the velocity (redshift space distortion) and the spacetime geometry (lensing).

The spectra depend sensitively and in several different ways on dark energy, on the matter and baryon densities, bias, etc. Their measurements provide a new estimation of cosmological parameters.

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