Cosmology with photometric redshift surveys: challenges and opportunities

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$2009 \rightarrow CERN$ summer student

$2012 \rightarrow PhD$

$2015 \rightarrow Postdoc$



 $2009 \rightarrow CERN su$ $2012 \rightarrow PhD$ $2015 \rightarrow Postdoc$



2019 \rightarrow Tenure track

Me and Geneva







(Un)observables **Observables Parameters** $A_s n_s$ $\Delta({f k},t)$ $f_{\rm NL}$ r $\Delta({f k},t)$ $\Delta^{lpha}(heta,\phi,\lambda)$ $\Omega_{\rm DE} \ \Omega_{\rm M}$ $w_a w_0$ Initial conditions Matter fluctuations $\alpha =:$ Energy components Power spectrum CMB temperature Background evolution CMB polarisation Galaxy density Galaxy shapes $Ly\alpha$ absorption 21cm flux





Photometric surveys

Galaxy clustering:

- $\delta_{g} = f[\delta_{M}] \sim b_{g} \delta_{M}$
- Local
- Spin-0





Photometric surveys

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- $e_i \sim \gamma_i \sim \delta_M$
- Spin-2



Photometric surveys

Galaxy clustering:

- $\delta_{g} = f[\delta_{M}] \sim b_{g} \delta_{M}$
- Local

Weak lensing:

• $e_i \sim \gamma_i \sim \delta_M$



Photometric surveys: the LSST



Outstanding numbers:

- World's largest imager
 8.4 m, 9.6 sq-deg FOV
- Wide: 20K sq-deg
- Deep: r~27
- Fast: ~100 visits per year
- Big data: ~15 TB per day

Dark Energy Science Collaboration:

- Supernovae
- Cluster science
- Strong lensing
- Weak lensing
- Large-scale structure



LSST Coll. et al. 0912.0201

Ideal analysis pipeline



- Photo-zs are complicated.
- Bunch galaxies up into photo-z bins and project onto the sphere.



DES Y1 data arXiv:1708.01530

2-point tomographic analysis

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- Compute all possible two-point crosscorrelations (different bins, different observables).



2-point tomographic analysis

- Photo-zs are complicated.
- Bunch galaxies up into photo-z bins and project onto the sphere.
- Compute all possible two-point crosscorrelations (different bins, different observables).
- Constrain parameters using a Gaussian likelihood.

 $-2 \log \mathsf{P}(\mathsf{d}|\theta) = (\mathsf{d}\text{-}\mathsf{t}(\theta))^{\mathsf{T}} \mathsf{C}^{-1} (\mathsf{d}\text{-}\mathsf{t}(\theta)) + \mathsf{L}_{0}$





Computing two-point functions



WIDE12H



A unified pseudo- C_{ℓ} estimator DA, F.J. Sanchez, A. Slosar arXiv:1809.09603





PCL facts

- Why C_{ℓ} ? (as opposed to $\xi(\theta)$)
 - k-cuts are easy to interpret. No Hankel transform
 - Covariance is a lot more diagonal
 - Good computational scaling (~N^{3/2})
- PCL vs. QMV
 - PCL == QMV when the covariance matrix is diagonal
 - PCL is precise enough in many common scenarios
 - QMV ~ N³, PCL ~ N^{3/2} (The trick is being able to estimate mode coupling analytically)



Efstathiou astro-ph/0307515 Leistedt et al. arXiv:1306.0005



http://www2.iap.fr/users/hivon/software/PolSpice/ https://gitlab.in2p3.fr/tristram/Xpol

What features does it implement?

- Calculate PCL power spectra (including coupling matrix, etc.)
- In curved and flat skies
- Spin-0 (density, CMB T) and spin-2 (shear, CMB Q/U) quantities
- Bells and whistles:
 - Mode deprojection
 - E/B mode purification
- Gaussian covariances



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Garcia-Garcia C., DA, Bellini E. arXiv:1906.11765





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<> Code (1) Issues 9 (1) Pull requests 3

LSSTDESC / NaMaster

Efstathiou astro-ph/0307515

A unified pseudo-Cl framework

A. Slosar: "The greatest thing since sliced bread"

- **Masking:** if I have a bad pixel, I make sure it doesn't get used.
- **Mode deprojection** is the extension of this idea into an arbitrary linear combination of pixels.

Imagine contaminating your data field as

Observed map
$$\rightarrow \delta^c_i = \delta^{-}_i + \alpha m_i$$
 Contaminant template (e.g. dust map)

True map

A proper analysis would marginalize over α .



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A proper analysis would marginalize over α .

If you do the maths, in PCL this amounts to:

- Finding the best fit value of α .
- Subtracting a contaminant map from the data using this α
- Calculate the PCL and correct for the bias this subtraction has produced
- Multiply by the inverse of the mode-coupling matrix

Leistedt et al. 1306.0005 Elsner et al. 1609.03577

NaMaster





Accuracy of $t(\theta) >>$ statistical power. LSST's statistical power will be awesome.

Requirements for LSST:

- Accuracy (errors well below statistical uncertainties)
- Robustness (thorough code validation and comparison)
- Flexibility (many observables, many cosmological models, ability to vary models and absorb systematics)
- Numerical performance (reasonable MCMC-ing time)

Core Cosmology Library: precision cosmological predictions for LSST Chisari E., DA, E. Krause +27, arXiv:1812.05995







Code: https://github.com/LSSTDESC/CCL Docs: https://ccl.readthedocs.io/en/latest/ Latest release: https://github.com/LSSTDESC/CCL/releases/tag/v2.0.1

Strict code validation requirements

- All calculations are performed with at least one different independent code.
- Agreement must be found within well-motivated/crazy stringent requirements.
- Alternative calculations are kept as benchmarks.
- CCL is automatically compared against benchmarks whenever a new addition is made to the code.
- Unit tested (~95%).

Code validation

Currently implemented:

- Background quantities and linear growth.
- Matter power spectrum Links to CAMB, CLASS, CosmicEmu, fast approximations
- Halo model:
 - Mass function
 - Bias
 - Concentrations
 - Profiles
 - Halo model power spectra
 - Easily generalisable
- Angular power spectra Galaxy clustering, cosmic shear, CMB lensing Easily generalisable
- Angular correlations functions
- 3D correlation functions

Used in a number of real-life analyses:

Cosmic shear: arXiv:1903.04957 Intrinsic alignments: arXiv:1911.01582,1901.09925 Cross-correlations: arXiv:1712.02738,1909.09102

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Covariances and data compression


Covariance matrices and data compression

A tomographic two-point function analysis already compresses the initial data vector significantly:

Catalogue with ~billions of objects and >5 quantities per object

A number of cross-correlations between sub-samples of these

What is the actual number of cross correlations?

Covariance matrices and data compression

A tomographic two-point function analysis already compresses the initial data vector significantly:

Catalogue with ~billions of objects and >5 quantities per object

A number of cross-correlations between sub-samples of these

What is the actual number of cross correlations?

Let's take an ideal LSST as an example:

10 redshift bins for lensing. 10 bins for clustering. 15 angular bins.

 $N_d = N_{\theta} N_{bin} (N_{bin}+1) / 2 = 3150$

Compression factor: $\sim 3 \times 10^6 \rightarrow$ pretty good!

Achieved by:

- Selecting only the most informative summary statistic.
- Averaging over equivalent modes (e.g. using statistical isotropy).

However, now we need to compute the data **covariance matrix**.

Different methods:

• Jackknife/bootstrap: use sub-samples of your own data.



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- Mock catalogues: based on N-body sims or fast methods (Gaussian, FLASK, 2LPT, COLA, PINOCHIO, PTHALOS, QuickPM ...)



Tassev et al. 1301.0322

For both of these, rule of thumb is $N_{samples} > 10 x$ (size of data vector). Then, O(3x10⁴) mocks/JKs are needed (covering the same volume as LSST).

Different methods:

- Jackknife/bootstrap: use sub-samples of your own data.
- Mock catalogues: based on N-body sims or fast methods (Gaussian, lognormal, 2LPT, COLA, PINOCHIO, PTHALOS, QuickPM ...)
- Analytical covariance matrix: Gaussian disconnected part:

 $Cov^{G} \left(C_{AB}^{ij}(l_{1}), C_{CD}^{kl}(l_{2}) \right) = \frac{4\pi \delta_{l_{1}l_{2}}}{\Omega_{s}(2l_{1}+1)\Delta l_{1}} \left[\left(C_{AC}^{ik}(l_{1}) + \delta_{ik}\delta_{AC}N_{A}^{i} \right) \left(C_{BD}^{jl}(l_{2}) + \delta_{jl}\delta_{BD}N_{B}^{j} \right) + \left(C_{AD}^{il}(l_{1}) + \delta_{il}\delta_{AD}N_{A}^{i} \right) \left(C_{BC}^{jk}(l_{2}) + \delta_{jk}\delta_{BC}N_{B}^{j} \right) \right]$

SSC

$$\operatorname{Cov}^{\operatorname{SSC}}\left(C_{AB}^{ij}(l_{1}), C_{CD}^{kl}(l_{2})\right) = \int d\chi \; \frac{q_{A}^{i}(\chi)q_{B}^{j}(\chi)q_{C}^{k}(\chi)q_{D}^{l}(\chi)}{\chi^{4}} \; \frac{\partial P_{AB}(l_{1}/\chi, z(\chi))}{\partial \delta_{b}} \frac{\partial P_{CD}(l_{2}/\chi, z(\chi))}{\partial \delta_{b}} \sigma_{b}(\Omega_{s}; z(\chi))$$

Relevant connected parts

$$\operatorname{Cov}^{\operatorname{NG},0}\left(C_{AB}^{ij}(l_1), C_{CD}^{kl}(l_2)\right) = \frac{1}{\Omega_{\mathrm{s}}} \int_{|\mathbf{l}| \in I_1} \frac{d^2 \mathbf{l}}{A(l_1)} \int_{|\mathbf{l}'| \in I_2} \frac{d^2 \mathbf{l}'}{A(l_2)} \int \mathrm{d}\chi \ \frac{q_A^i(\chi) q_B^j(\chi) q_C^k(\chi) q_D^l(\chi)}{\chi^6} \ T_{ABCD}^{ijkl}(\mathbf{l}/\chi, -\mathbf{l}/\chi, \mathbf{l}'/\chi, -\mathbf{l}'/\chi; z(\chi))$$

- + double Hankel transform if you work in real space
- + probably worry about survey geometry (mode coupling)

$$\left\langle \Delta \tilde{C}_{\ell}^{ab} \Delta \tilde{C}_{\ell'}^{cd} \right\rangle = \sum_{mm'} \sum_{\mathbf{l}_1 \mathbf{l}_2} \left(C^{ac}_{\ell_1} C^{bd}_{\ell_2} W^a_{\mathbf{l}\mathbf{l}_1} W^b_{\mathbf{l}\mathbf{l}_2} W^c_{\mathbf{l}\mathbf{l}_1} W^d_{\mathbf{l}\mathbf{l}_2} + C^{ad}_{\ell_1} C^{bc}_{\ell_2} W^a_{\mathbf{l}\mathbf{l}_1} W^b_{\mathbf{l}\mathbf{l}_2} W^d_{\mathbf{l}\mathbf{l}_1} \right)$$

Computation scales very bad: $O(N_{\theta}^2 N_{bin}^4)$

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Krause & Eifler 1601.05779

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Garcia-Garcia et al. arXiv:1906.11765

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All of these cases would benefit massively from reducing the size of the data vector.

Can we compress further?

Science-driven 3D data compression DA, arXiv:1707.08950

Sheer shear: weak lensing with one mode E. Bellini, DA et al. arXiv:1903.04957



Idea: find the linear combinations of your data that contain most of the information about a given parameter θ .

$$y_p \equiv \mathbf{e}_p^\dagger \mathbf{x}$$

Data: $\mathbf{x} \rightarrow \text{maps}/a_{\ell m}$ s of a given set of tomographic observables

(e.g. galaxy overdensity or shear in a set of redshift bins).

The linear coefficients **e** can be found as the eigenvectors of a generalized eigenvalue equation:

$$\partial_{\theta} \mathsf{C} \mathbf{e}_p = \lambda_p \mathbf{C} \mathbf{e}_p$$

One generic parameter we could optimize for is the overall S/N amplitude. Maximizing this should provide us with most of the information about any parameter in most cases.

In this case, the eigenvalue equation reads: Signal covariance $(S + N)e_p = \lambda_p Ne_p$

Resulting modes y_p are uncorrelated and contain the maximum amount of information (info(y_0) > info(y_1) > ...).

Example: galaxy clustering with spectroscopic redshifts.

- $\mathbf{x} \rightarrow$ galaxy overdensity in an infinitesimal redshift bin.
- $C \rightarrow$ all possible cross-power spectra between bins (noise + signal) $N \rightarrow$ flat, diagonal shot-noise power spectrum

The solution to the generalized eigenvalue equation (KL modes) is

$$e_{k,\ell}(z) \propto j_{\ell} \left(k \chi(z) \right)$$

i.e. KL transform in this case is the harmonic-Bessel transform.

The covariance of the resulting KL modes is

$$\lambda_{k,\ell} \propto P(k)$$

i.e. in this case the KL transform tells you to just compute the Fourier transform and estimate the 3D power spectrum (as expected!).

The KL eigenmodes are the generalization of a P(k) analysis to other types of data.

Data compression



 χ [Mpc]

Data compression



You **can** compress further!



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- Majority of the signal concentrated in 1st KL mode.



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- Size of data vector: $2 \times 5 \times (7 \times 8)/2 = 280$ elements.
- Eigenvectors close to scale-independent. Think of them as redshift-dependent galaxy weights.
- The first 1-2 modes are able to recover the full constraining power. **Compression factor ~19-30**!

Other uses of the KL transform:

- Large-scale effects: optimize fNL constraints.
- Systematics: remove modes that are most sensitive to e.g. intrinsic alignments, magnification ...
 (basically, put even thing you don't like in the poice component)

(basically put everything you don't like in the noise component)

• Foreground removal in 21cm experiments

Extreme data compression:

- Alsing & Wandelt 1712.00012, Alsing et al. 1801.01497.
- One summary statistic per free parameter.
- Can be made robust to systematics.
- Potentially more sensitive to modeling errors. Missing systematics may be more difficult to detect (KL at least gives you maps to inspect).

Gaussian likelihood

$$-2 \log P(\mathbf{d}|\theta) = (\mathbf{d} - \mathbf{t}(\theta))^{\mathsf{T}} \mathbf{C}^{-1}(\theta) (\mathbf{d} - \mathbf{t}(\theta)) + \mathbf{L}_0 ?$$

- Do we have to take into account the parameter dependence of the covariance matrix?
- I.e. do we need to compute a new covariance at every point in an MCMC chain?

The effect on cosmological parameter estimation of a parameter-dependent covariance matrix Kodwani D., DA, P. Ferreira

arXiv:1811.11584





-2 log P(d| θ) = (d-t(θ))^T C⁻¹(θ) (d-t(θ)) + L₀ ?

- Do we have to take into account the parameter dependence of the covariance matrix?
- I.e. do we need to compute a new covariance at every point in an MCMC chain?
- Carron 2016: <u>for Gaussian</u> fields it's not only unnecessary, <u>it's</u> <u>incorrect</u>.
- The galaxy overdensity and cosmic shear aren't Gaussian, so do we need to worry about this at all?

The math

The information content of the covariance matrix can be quantified approximating the likelihood as Gaussian around the maximum (i.e. a la Fisher).

• Effect on parameter uncertainties:

$$\mathcal{F}_{\mu\nu} = \partial_{\mu} \mathbf{t}^T \Sigma^{-1} \partial_{\nu} \mathbf{t} + \frac{1}{2} \operatorname{Tr} \left(\Sigma^{-1} \partial_{\mu} \Sigma \Sigma^{-1} \partial_{\nu} \Sigma \right)$$

• Effect on parameter bias:

$$\Delta \theta_{\mu} = -\frac{1}{2} \mathcal{F}_{\mu\nu}^{-1} \, \mathcal{F}_{\rho\tau}^{-1} \, \partial_{\rho} \mathbf{t}^{T} \boldsymbol{\Sigma}^{-1} \partial_{\nu} \boldsymbol{\Sigma} \, \boldsymbol{\Sigma}^{-1} \partial_{\tau} \mathbf{t}$$

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Let's examine the dependence on f_{sky} .

Roughly: $\Sigma \propto f_{\rm sky}^{-1}$ Then: $\Delta \theta \propto f_{\rm sky}^{-1}$, $\delta \sigma(\theta) \propto f_{\rm sky}^{-3/2}$

In general, the effects of a parameter-dependent covariance shrink with the number of modes in the analysis (same also with ℓ_{max}).

Parameter-dependent covariances

Results: parameter uncertainties



The parameter dependence of the covariance is irrelevant in all cases.

Tomographic galaxy clustering with the Subaru Hyper Suprime-Cam first-year public data release Nicola A., DA, Slosar A., F.J. Sanchez

NICOIA A., DA, SIOSAR A., F.J. Sanchez et al. (LSST DESC) (arXiv coming soon!)







The HSC survey



- HSC end goal: 5-year survey covering 1400 sq-deg
- Deep ($r_{lim} \sim 26$), very good seeing (0.6")
- 5 bands: grizy
- 1st-year data release: full depth on ~150 sq-deg in 6 fields (+ a few deeper fields)
- Precursor to LSST. Common DM pipeline, similar depth.

The HSC survey



Hikage et al. arXiv:1809.09148

Cosmological constraints from shear

HSC & LSST

Why is DESC interested?

- Same DM pipeline, similar depth: ideal testbed for our pipelines.
- Test viability of Fourier-space clustering analysis in the presence of sky systematics.
- Photometric clustering has focused on small samples with good photo-zs (e.g. LRGs, redMaGiC)
 - This will mean losing almost all of our galaxies in LSST. Can we do better?
 - How much more stringent are photo-z calibration requirements for galaxy clustering?
- No-one has looked at photometric clustering in HSC.



Cut	Comment
<pre>detect_is_primary=True icmodel_flags_badcentroid=False icentroid_sdss_flags=False iflags_pixel_edge=False iflags_pixel_interpolated_center=False iflags_pixel_saturated_center=False iflags_pixel_cr_center=False iflags_pixel_bad=False iflags_pixel_suspect_center=False iflags_pixel_clipped_any=False meas.ideblend_skipped=False</pre>	Basic quality cuts, see [35, 43]
<pre>[g,r,z,y]centroid_sdss_flags=False [g,r,i,z,y]cmodel_flux_flags=False [g,r,i,z,y]flux_psf_flags=False [g,r,z,y]flags_pixel_edge=False [g,r,z,y]flags_pixel_interpolated_center=False [g,r,z,y]flags_pixel_saturated_center=False [g,r,z,y]flags_pixel_cr_center=False [g,r,z,y]flags_pixel_bad=False</pre>	Strict photometry cuts
icmodel_mag-a_i < 24.5	Magnitude limit
icmodel_flux> 10 icmodel_flux_err	10σ detections
[g,r,y,z]cmodel_flux>5[g,r,y,z]cmodel_flux_err	5σ detection (required only in 2 other bands)
iclassification_extendedness=1	Star-galaxy separator

- Similar to HSC shear sample (no shape cuts)
- Bright magnitude cut (5σ limit is ~26)
- This improves sample homogeneity



- Similar to HSC shear sample (no shape cuts)
- Bright magnitude cut (5σ limit is ~26)
- This improves sample homogeneity
- Split into 4 redshift bins
- Photo-z posteriors from several codes



- We pixelise each field using square **pixels 0.6'** in size.
- Plate-Carrée projection.
- Small fields (<20 sq-deg) \rightarrow **flat-sky** approximation.
- Mask: footprint + bright-object mask + depth mask.
- Overdensity maps for each field and redshift bin.

$$\delta_p = \frac{n_p}{\bar{n}w_p} - 1 \qquad \bar{n} = \frac{\langle n_p w_p \rangle}{\langle w_p \rangle}$$



- We generate sky maps for all quantities that could potentially cause systematic fluctuations in δ_g .
- Observing conditions are mapped in all bands.
- 47 maps in total (per field).
- We deproject all of these in all power spectra.

Redshift distributions

• The redshift distributions are a central part of the theory prediction

$$C_{\ell}^{ab} = \int d\chi \, \frac{H^2}{\chi^2} \, p_a(z) \, p_b(z) \, P_{ab}\left(k = \ell + \frac{1}{2}, z\right)$$

- Our fiducial distributions are determined from the COSMOS 30band photometric catalog.
- COSMOS objects are reweighted in color space to match our sample.
- We obtain alternative estimates of the p(z)s by stacking the pdfs of all objects in each bin for 4 different photo-z codes.



Redshift distributions

- None of these estimates are exact: we need to marginalise over residual uncertainties.
- N(z) uncertainties have traditionally been summarized by a single "shift" parameter in shear analyses.
- Clustering is potentially more sensitive to this, so we extend this by adding a "width" parameter.
- We vary these within broad priors.

$$p_i(z) = \hat{p}_i(z_c + (1 + z_{w,i})(z - z_c) + \Delta z_i)$$



Power spectra





- Power spectra estimated with NaMaster
- 47 templates deprojected.
- Analytic shot-noise correction.

(3, 3)

(3, 4)

 10^{4}



We use an **analytical covariance** matrix that includes:

• Mode coupling in the Gaussian part due to **survey geometry** (computed with NaMaster).

$$\operatorname{Cov}\left(\tilde{C}_{\ell}^{ab}, \tilde{C}_{\ell'}^{cd}\right) = (2\ell'+1)^{-1} \left[C_{(\ell}^{ac} C_{\ell'}^{bd} M_{\ell\ell'}(w_a w_c, w_b w_d) + C_{(\ell}^{ad} C_{\ell'}^{bc} M_{\ell\ell'}(w_a w_d, w_b w_c) \right]$$

 Mode-coupling due to non-linear growth using a perturbation theory + halo model approach.

$$\begin{aligned} \operatorname{Cov}_{\mathrm{NG}}(C_{\ell}^{ab}, C_{\ell'}^{cd}) &= \frac{1}{4\pi f_{\mathrm{sky}}} \int_{|\boldsymbol{\ell}| \in \ell_1} \int_{|\boldsymbol{\ell}'| \in \ell_2} \int \frac{\mathrm{d}^2 \boldsymbol{\ell}}{A(\ell_1)} \frac{\mathrm{d}^2 \boldsymbol{\ell}'}{A(\ell_2)} \,\mathrm{d}\chi \, \frac{q^a(\chi) q^b(\chi) q^c(\chi) q^d(\chi)}{\chi^6} \times \\ T^{abcd}(\boldsymbol{\ell}/\chi, -\boldsymbol{\ell}/\chi, \boldsymbol{\ell}'/\chi, -\boldsymbol{\ell}'/\chi). \end{aligned}$$
$$T^{abcd} &= T^{abcd, 1h} + (T_{22}^{abcd, 2h} + T_{13}^{abcd, 2h}) + T^{abcd, 3h} + T^{abcd, 4h} \end{aligned}$$

• Mode-coupling due to **super-survey** modes.

$$\begin{aligned} \operatorname{Cov}_{\mathrm{SSC}}(C_{\ell}^{ab}, C_{\ell'}^{cd}) &= \int \mathrm{d}\chi \; \frac{q^{a}(\chi)q^{b}(\chi)q^{c}(\chi)q^{d}(\chi)}{\chi^{4}} \times \\ & \frac{\partial P_{ab}(\ell/\chi, z(\chi))}{\partial \delta_{\mathrm{LS}}} \frac{\partial P_{cd}(\ell'/\chi, z(\chi))}{\partial \delta_{\mathrm{LS}}} \sigma_{b}^{2}(z(\chi)) \end{aligned}$$

• Covariances are estimated in each field and then coadded.
Covariances are model dependent!

We use a four-step process:

- 1. Gaussian covariance from measured power spectra.
- 2. Obtain best-fit parameters and compute corresponding covariance.
- 3. Run chains with this covariance.
- 4. If new best-fit is too far from the previous one, GOTO 2.

Covariance matrices



Robustness tests: consistency across fields



Robustness tests: masks and contaminants



Robustness tests



HOD parameterisation

$$P_{gg}(z,k) = P_{gg,1h}(z,k) + P_{gg,2h}(z,k)$$



$$P_{gg,1h}(k) = \frac{1}{\bar{n}_a^2} \int dM \, \frac{dn}{dM} \bar{N}_c \left[\bar{N}_s^2 u_s^2(k) + 2\bar{N}_s u_s^2(k) \right]$$
$$P_{gg,2h}(k) = \left(\frac{1}{\bar{n}_g} \int dM \, \frac{dn}{dM} \, b_h(M) \, \bar{N}_c \left[1 + \bar{N}_s u_s(k) \right] \right)^2 P_{\text{lin}}(k)$$

HOD parameterisation













Results: Magnification



Results: Magnification







Results: photo-z systematics



Results: galaxy bias



Summary

- Tools for cosmological analyses:
 - Power spectra and Gaussian covariances: NaMaster (arXiv:1809.09603).
 - Accurate theory predictions: CCL (arXiv:1812.05995).
- Data compression for weak lensing (arXiv:1903.04957):
 - KL transform == P(k) analysis for weak lensing.
 - Information contained in a small number of modes N.
 - N=1-2 for current data.
- Parameter-dependent covariances (arXiv:1811.11584):
 - Don't worry about this.
- Galaxy clustering in HSC DR1 (coming soon):
 - Fourier-space analysis with comprehensive systematic deprojection.
 - Mild redshift dependence of HOD due to red galaxy dropout.
 - Simple prescription for b(z, m_{lim})
 - Additional sensitivity to photo-z uncertainties (challenge and opportunity)
 - 3σ detection of cosmic magnification
 - Consistent cosmological constraints.

Merci beaucoup!